



Improved energy bounds for Schrödinger operators

Lorenzo Brasco, Giuseppe Buttazzo

► To cite this version:

Lorenzo Brasco, Giuseppe Buttazzo. Improved energy bounds for Schrödinger operators. *Calculus of Variations and Partial Differential Equations*, 2015, 53, pp.977-1014. 10.1007/s00526-014-0774-1 . hal-00955139v2

HAL Id: hal-00955139

<https://hal.science/hal-00955139v2>

Submitted on 28 Jul 2014

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Copyright

IMPROVED ENERGY BOUNDS FOR SCHRÖDINGER OPERATORS

LORENZO BRASCO AND GIUSEPPE BUTTAZZO

ABSTRACT. Given a potential V and the associated Schrödinger operator $-\Delta + V$, we consider the problem of providing sharp upper and lower bound on the energy of the operator. It is known that if for example V or V^{-1} enjoys suitable summability properties, the problem has a positive answer. In this paper we show that the corresponding isoperimetric-like inequalities can be improved by means of quantitative stability estimates.

CONTENTS

1. Introduction	1
1.1. Aim of the paper	1
1.2. Main results	3
1.3. Plan of the paper	4
2. Preliminaries	4
3. Maximization problems	8
4. Stability for maximization problems	9
4.1. Stability of the potentials	10
4.2. Stability of the state functions	12
5. Minimization problems	13
6. Stability for minimization problems	15
6.1. Preliminary results	15
6.2. Stability of the potentials	17
6.3. Stability of the state functions	24
Appendix A. Sharp decay estimates for non autonomous Schrödinger equations	24
References	30

1. INTRODUCTION

1.1. Aim of the paper. Let $\Omega \subset \mathbb{R}^N$ be an open set not necessarily with finite measure (it could be $\Omega = \mathbb{R}^N$) and $V : \Omega \rightarrow \mathbb{R}$ be a potential. We consider the associated Schrödinger operator $-\Delta + V$ defined on the homogeneous Sobolev space $W_0^{1,2}(\Omega)$. The latter is the closure of $C_0^\infty(\Omega)$ with respect to the norm

$$\|u\|_{W_0^{1,2}(\Omega)} := \left(\int_{\Omega} |\nabla u|^2 dx \right)^{1/2}.$$

We also denote by $W^{-1,2}(\Omega)$ its dual space and by $\langle \cdot, \cdot \rangle$ the duality pairing between $W_0^{1,2}(\Omega)$ and $W^{-1,2}(\Omega)$. In this paper we are concerned with the following problem: given a source term $f \in W^{-1,2}(\Omega)$, find lower and upper bounds on the *energy* of the relevant Schrödinger operator, i.e.

$$\mathcal{E}_f(V) = -\frac{1}{2} \int_{\Omega} |\nabla u_V|^2 dx - \frac{1}{2} \int_{\Omega} V u_V^2 dx.$$

Date: 14th July 2014, Marseille et Pise.

2010 Mathematics Subject Classification. 49K20, 49Q10, 49Q20.

Key words and phrases. Schrödinger operators; optimal potentials; stability inequalities; decay estimates.

Here the state function u_V is a $W_0^{1,2}(\Omega)$ solution of

$$-\Delta u + V u = f, \quad \text{in } \Omega.$$

In the recent paper [10], this problem has been solved for summable potentials or for *confining potentials*, i.e. for potentials blowing-up at infinity such that $1/V$ enjoys some summability properties. For example, the harmonic-like potential

$$V = (\delta^2 + |x|^2)^{\gamma/2}, \quad \delta > 0,$$

belongs to this class, for suitable $\gamma > 0$. In order to provide a deeper insight into the scopes of this work, it is useful to briefly recall some of the results in [10]. In that paper it has been shown that $\mathcal{E}_f(V)$ can be universally bounded *from above* in the class (see [10, Proposition 5.1])

$$\mathcal{V}_1 = \left\{ V : \int_{\Omega} |V|^p \leq 1 \right\},$$

and *from below* in the class (see [10, Proposition 5.4])

$$\mathcal{V}_2 = \left\{ V \geq 0 : \int_{\Omega} V^{-p} \leq 1 \right\}.$$

The value 1 above plays no special role and can be replaced by any constant $c > 0$. Indeed, one can show that there exist two potentials V_0 and U_0 such that

$$\int_{\Omega} |V_0|^p dx = 1 = \int_{\Omega} |U_0|^{-p} dx,$$

and

$$(1.1) \quad \mathcal{E}_f(V) \leq \mathcal{E}_f(V_0), \quad \text{for every } V \in \mathcal{V}_1,$$

$$(1.2) \quad \mathcal{E}_f(V) \geq \mathcal{E}_f(U_0), \quad \text{for every } V \in \mathcal{V}_2.$$

In both the estimates (1.1) and (1.2), the equality sign holds if and only if $V = V_0$ and $V = U_0$, respectively. Moreover, the extremal potentials V_0 and U_0 can be characterized in terms of the solutions v_0 and u_0 of the semilinear non-autonomous PDEs

$$-\Delta v_0 + c v_0^{(p+1)/(p-1)} = f \quad \text{and} \quad -\Delta u_0 + c u_0^{(p-1)/(p+1)} = f,$$

through the relations

$$V_0 = \left(\int_{\Omega} |v_0|^{2p/(p-1)} dx \right)^{-1/p} |v_0|^{2/(p-1)} \quad \text{and} \quad U_0 = \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{1/p} |u_0|^{-2/(p+1)}.$$

For completeness, we point out that a special class of potentials from the sets \mathcal{V}_1 and \mathcal{V}_2 are given respectively by

$$(1.3) \quad V(x) = \begin{cases} |E|^{-1/p}, & x \in E, \\ 0, & \text{otherwise} \end{cases} \quad \text{and} \quad V(x) = \begin{cases} |E|^{1/p}, & x \in E, \\ +\infty, & \text{otherwise,} \end{cases}$$

where $E \subset \Omega$ is an open set and $|E|$ denotes its N -dimensional Lebesgue measure. If we suppose for simplicity that $|E| = 1$, the corresponding operators are given by

$$W_0^{1,2}(\Omega) \ni u \mapsto -\Delta u + u \cdot 1_E \quad \text{and} \quad W_0^{1,2}(E) \ni u \mapsto -\Delta u + u.$$

Thus from (1.1) and (1.2) we get

$$\min_{u \in W_0^{1,2}(\Omega)} \left[\frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2} \int_E |u|^2 dx - \langle f, u \rangle \right] < \mathcal{E}_f(V_0),$$

and

$$\min_{u \in W_0^{1,2}(E)} \left[\frac{1}{2} \int_E |\nabla u|^2 dx + \frac{1}{2} \int_E |u|^2 dx - \langle f, u \rangle \right] > \mathcal{E}_f(U_0),$$

for every $E \subset \Omega$ with $|E| = 1$. Observe that for the second problem the set Ω simply acts as a *design region* where the admissible domains E have to be contained.

The problem of finding sharp bounds on energetical quantities linked to a Schrödinger operator is quite classical, with many studies devoted to the *ground state energy* or *first eigenvalue*

$$\lambda_1(V) = \min_{u \in W_0^{1,2}(\Omega)} \left\{ \int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} V u^2 dx : \|u\|_{L^2(\Omega)} = 1 \right\}.$$

For example, the pioneering paper [22] by Keller considers the problem of finding sharp lower bounds on $\lambda_1(V)$ in the class \mathcal{V}_1 , in the case of space dimension $N = 1$ and $\Omega = \mathbb{R}$. Related problems have been considered by Ashbaugh and Harrell in [2] in higher dimensions. There is a vast literature on the subject, considering optimal bounds for other spectral quantities, like the *first excited state* or *second eigenvalue* $\lambda_2(V)$ and the *fundamental gap* $\lambda_2(V) - \lambda_1(V)$. We also mention the recent paper [8], where the case of successive eigenvalues $\lambda_k(V)$ for $k \geq 2$ is considered, for the non-compact case of $\Omega = \mathbb{R}^N$. Actually this kind of problems is even older: indeed, we observe that for potentials of the second form in (1.3), we have

$$\lambda_1(V) = \min_{u \in W_0^{1,2}(E)} \left\{ \int_E |\nabla u|^2 dx + \int_E |u|^2 dx : \|u\|_{L^2(E)} = 1 \right\} = \lambda_1(E) + 1,$$

where $\lambda_1(E)$ now stands for the first eigenvalue of $-\Delta$ with Dirichlet boundary conditions on ∂E . Thus the celebrated *Faber-Krahn inequality* (see [21, Chapter 3])

$$\lambda_1(E) \geq \lambda_1(B), \quad \text{with } B \text{ any ball such that } |B| = |E| = 1,$$

can be seen as a particular instance of these problems. A more general overview on optimization problems of spectral type can be found in [9] and [21], to which we refer the interested reader. We wish to point out that on the contrary the case of the energy $\mathcal{E}_f(V)$ appears to be less investigated.

1.2. Main results. In this paper, we improve the previous sharp bounds (1.1) and (1.2) on the energy of a Schrödinger operator, by means of a quantitative stability result. In other words, we will prove that the energy gap $\mathcal{E}_f(V_0) - \mathcal{E}_f(V)$ or $\mathcal{E}_f(V) - \mathcal{E}_f(U_0)$ controls the deviation from optimality of a potential V . Thus it is possible to add a reminder term in the right-hand side of (1.1) and (1.2), which measures the distance of a generic potential V from V_0 or U_0 . The relevant results are summarized in the following couple of theorems, which represent the main results of the paper. We refer the reader to Sections 4 and 6 for the precise statements.

Theorem A (Stability of the maximizer). *Let $1 < p < \infty$. There exists a constant $\sigma_1 > 0$ such that for every $V \in \mathcal{V}_1$ we have*

$$\begin{cases} \mathcal{E}_f(V) \leq \mathcal{E}_f(V_0) - \sigma_1 \|V - V_0\|_{L^p(\Omega)}^2, & \text{if } p \geq 2, \\ \mathcal{E}_f(V) \leq \mathcal{E}_f(V_0) - \sigma_1 \left\| |V|^{p-2} V - |V_0|^{p-2} V_0 \right\|_{L^{p'}(\Omega)}^2, & \text{if } 1 < p < 2. \end{cases}$$

In the case of inequality (1.2) we need to distinguish between the case $|\Omega| < +\infty$ and $|\Omega| = +\infty$.

Theorem B (Stability of the minimizer). *Let $\Omega \subset \mathbb{R}^N$ be an open set such that $|\Omega| = +\infty$. Let $1 < p < \infty$, $r > N/2$, and let $f \in L^r(\Omega)$ be a function decaying to 0 at infinity as $O(|x|^{-\alpha})$ with $\alpha > 1 + N/2$. Then there exist a constant $\sigma_2 > 0$ and an exponent $\beta = \beta(p) > 2$ such that for every $V \in \mathcal{V}_2$ we have*

$$\mathcal{E}_f(V) \geq \mathcal{E}_f(U_0) + \sigma_2 \left\| \frac{1}{V} - \frac{1}{U_0} \right\|_{L^p(\Omega)}^\beta.$$

If $|\Omega| < +\infty$, the same result holds for every $f \in W^{-1,2}(\Omega)$ without any additional hypothesis.

Stability results of this type have attracted an increasing interest in recent years. As a non-exhaustive list of works on the subject, we point out for example [13] and [17] dealing with the classical isoperimetric inequality, the papers [5, 6, 7, 20] and [23] concerning sharp bounds

for eigenvalues of the Laplacian and [4, 11, 15] about quantitative versions of the Sobolev and Gagliardo-Nirenberg inequalities with sharp constant.

Among these papers, the recent one [12] is very much related with the subject here considered. In [12] a quantitative stability estimate for $\lambda_1(V)$ is proved, for potentials belonging to the class (here $r > N/2$)

$$\mathcal{V}' = \left\{ V : \int_{\mathbb{R}^N} |V_-|^r \leq 1 \right\},$$

where V_- is the negative part of V . In this case $\lambda_1(V)$ admits a sharp lower bound, corresponding to the negative potential

$$W_0 = - \left(\int_{\mathbb{R}^N} |w_0|^{2r/(r-1)} \right)^{-1/r} |w_0|^{2r/(r-1)},$$

where w_0 is an extremal of the Gagliardo-Nirenberg inequality

$$\left(\int_{\mathbb{R}^N} |u|^{2r/(r-1)} dx \right)^{(r-1)/r} \leq C \left(\int_{\mathbb{R}^N} |\nabla u|^2 dx \right)^{1-\vartheta} \left(\int_{\mathbb{R}^N} |u|^2 dx \right)^{\vartheta}.$$

The parameter $0 < \vartheta = \vartheta(N, r) < 1$ above is uniquely determined by scaling invariance.

1.3. Plan of the paper. Along all the paper, for the sake of simplicity, we assume $N \geq 3$. In this way the Sobolev exponent 2^* is finite; all the results also apply, with minor modifications, to the cases $N = 1$ and $N = 2$ by using the corresponding Sobolev embeddings.

We start with Section 2 where we fix the main notations and prove some basic results which will be used throughout the whole paper. Then Section 3 is concerned with maximization problems for \mathcal{E}_f , under a constraint on the L^p norm of the admissible potentials. The relevant quantitative stability result (Theorem 4.3) is then considered in Section 4. Minimization problems are addressed in Section 5, while the last section of the paper contains the corresponding stability result (Theorem 6.6). Finally, a self-contained Appendix on sharp decay estimates for finite energy solutions of

$$-\Delta u + c u^{q-1} = f, \quad \text{for } c > 0, \quad 1 < q < 2,$$

concludes the paper (Theorem A.1).

2. PRELIMINARIES

In the paper we mainly focus on the following three model cases:

- $\Omega = \mathbb{R}^N$;
- $\Omega = \omega \times \mathbb{R}$, with $\omega \subset \mathbb{R}^{N-1}$ open set with finite Lebesgue measure (*waveguide*);
- $\Omega \subset \mathbb{R}^N$ with finite Lebesgue measure (*compact case*).

For $N \geq 3$, we define

$$2^* = \frac{2N}{N-2}.$$

The following embedding properties of $W_0^{1,2}(\Omega)$ are well known.

Proposition 2.1. *Let $N \geq 3$; then*

- i) *if $\Omega = \mathbb{R}^N$, we have the continuous embedding $W_0^{1,2}(\Omega) \hookrightarrow L^{2^*}(\Omega)$, but $W_0^{1,2}(\Omega) \not\hookrightarrow L^s(\Omega)$ for $s \neq 2^*$;*
- ii) *if $\Omega = \omega \times \mathbb{R}$ is a waveguide, we have the continuous embedding $W_0^{1,2}(\Omega) \hookrightarrow L^s(\Omega)$ for every $2 \leq s \leq 2^*$;*
- iii) *if $|\Omega| < +\infty$, we have the continuous embedding $W_0^{1,2}(\Omega) \hookrightarrow L^s(\Omega)$ for every $0 < s \leq 2^*$. Moreover, this is compact for $0 < s < 2^*$.*

In what follows, for $N \geq 3$ we set

$$(2.1) \quad T_N := \inf_{u \in W_0^{1,2}(\mathbb{R}^N)} \left\{ \int_{\mathbb{R}^N} |\nabla v|^2 dx : \|v\|_{L^{2^*}(\mathbb{R}^N)} = 1 \right\} < +\infty.$$

This infimum is finite by Proposition 2.1 and attained on $W_0^{1,2}(\mathbb{R}^N)$, see for example [26].

Let $f \in W^{-1,2}(\Omega)$; for every potential V belonging to the admissible class

$$\mathcal{V} = \left\{ V : \Omega \rightarrow (-\infty, +\infty] : V \text{ Borel measurable, } \|V_-\|_{L^{N/2}(\Omega)} < T_N \right\},$$

we define its energy by

$$(2.2) \quad \mathcal{E}_f(V) = \min_{u \in W_0^{1,2}(\Omega)} \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2} \int_{\Omega} V u^2 dx - \langle f, u \rangle.$$

Proposition 2.2. *For every $V \in \mathcal{V}$ the minimization problem (2.2) admits a solution $u_V \in W_0^{1,2}(\Omega)$. Moreover the energy inequality*

$$(2.3) \quad \|u_V\|_{W_0^{1,2}(\Omega)} \leq \frac{T_N}{T_N - \|V_-\|_{L^{N/2}(\Omega)}} \|f\|_{W^{-1,2}(\Omega)}.$$

holds.

Proof. At first we observe that, taking $u = 0$ gives $\mathcal{E}_f(V) \leq 0$. Moreover, since $V \in \mathcal{V}$, the energy functional is bounded from below, because

$$(2.4) \quad \int_{\Omega} V u^2 dx \geq - \int_{\Omega} V_- u^2 dx \geq - \|V_-\|_{L^{N/2}(\Omega)} \|u\|_{L^{2^*}(\Omega)}^2 \geq - \frac{\|V_-\|_{L^{N/2}(\Omega)}}{T_N} \int_{\Omega} |\nabla u|^2 dx,$$

and thus

$$(2.5) \quad \begin{aligned} & \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2} \int_{\Omega} V u^2 dx - \langle f, u \rangle \\ & \geq \frac{1}{2} \left(1 - \frac{\|V_-\|_{L^{N/2}(\Omega)}}{T_N} - \delta \right) \int_{\Omega} |\nabla u|^2 dx - \frac{1}{2\delta} \|f\|_{W^{-1,2}(\Omega)}^2, \end{aligned}$$

for every $0 < \delta \leq 1$. Thus $\mathcal{E}_f(V)$ is finite. Let now $\{u_n\}_{n \in \mathbb{N}} \subset W_0^{1,2}(\Omega)$ be a minimizing sequence; we can assume that

$$\frac{1}{2} \int_{\Omega} |\nabla u_n|^2 dx + \frac{1}{2} \int_{\Omega} V u_n^2 dx - \langle f, u_n \rangle \leq \mathcal{E}_f(V) + 1.$$

By (2.5), if we take $\delta \ll 1$ the sequence $\{u_n\}_{n \in \mathbb{N}}$ is bounded in $W_0^{1,2}(\Omega)$, so u_n weakly converges (up to a subsequence) in $W_0^{1,2}(\Omega)$ to a function $u \in W_0^{1,2}(\Omega)$. Moreover, by the compact Sobolev embedding of Proposition 2.1 iii), we have strong convergence in $L^s(\Omega')$, for every (smooth) $\Omega' \Subset \Omega$ and every $1 \leq s < 2^*$. In particular, u_n converges almost everywhere (up to a subsequence) in Ω to u . Also observe that still by Proposition 2.1 we have weak convergence of u_n^2 to u^2 in $L^{2^*/2}(\Omega)$. By using this and the Fatou Lemma, we get

$$\liminf_{n \rightarrow \infty} \int_{\Omega} V u_n^2 dx = \liminf_{n \rightarrow \infty} \left[\int_{\Omega} V_+ u_n^2 dx - \int_{\Omega} V_- u_n^2 dx \right] \geq \int_{\Omega} V u^2 dx.$$

Finally, the weak lower semicontinuity of the norm implies

$$\begin{aligned} & \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2} \int_{\Omega} V u^2 dx - \langle f, u \rangle \\ & \leq \liminf_{n \rightarrow \infty} \left[\frac{1}{2} \int_{\Omega} |\nabla u_n|^2 dx + \frac{1}{2} \int_{\Omega} V u_n^2 dx - \langle f, u_n \rangle \right] = \mathcal{E}_f(V), \end{aligned}$$

which gives the existence of a minimizer u_V .

This function u_V satisfies the Euler-Lagrange equation

$$\int_{\Omega} \langle \nabla u, \nabla \varphi \rangle dx + \int_{\Omega} V u \varphi dx = \langle f, \varphi \rangle,$$

for every $\varphi \in C_0^\infty(\Omega) \cap L^2(\Omega; V)$, where for $V \in \mathcal{V}$ we set

$$L^2(\Omega; V) = \left\{ \varphi : \int_{\Omega} |V| \varphi^2 dx < +\infty \right\}.$$

By density, the previous equation holds for every $\varphi \in W_0^{1,2}(\Omega) \cap L^2(\Omega; V)$. By taking u_V as a test function and then appealing to

$$|\langle f, u \rangle| \leq \frac{1}{2\delta} \|f\|_{W^{-1,2}(\Omega)}^2 + \frac{\delta}{2} \|u\|_{W_0^{1,2}(\Omega)}^2,$$

we have the estimate

$$\left(1 - \frac{\delta}{2}\right) \int_{\Omega} |\nabla u_V|^2 dx + \int_{\Omega} V u_V^2 dx \leq \frac{1}{2\delta} \|f\|_{W^{-1,2}(\Omega)}^2.$$

By using (2.4) and choosing

$$\delta = 1 - \frac{\|V_-\|_{L^{N/2}(\Omega)}}{T_N},$$

we get (2.3). □

Remark 2.3. If $\|V_-\|_{L^{N/2}} \geq T_N$, in general problem (2.2) is not well-posed. Indeed, take $\Omega = \mathbb{R}^N$ and indicate by $U \in W_0^{1,2}(\mathbb{R}^N)$ a positive function such that

$$\int_{\mathbb{R}^N} |\nabla U|^2 dx = T_N \quad \text{and} \quad \int_{\mathbb{R}^N} |U|^{2^*} dx = 1.$$

We then take

$$V = -T_N U^{4/(N-2)},$$

and $f \in L^{(2^*)}'(\mathbb{R}^N)$ such that

$$\int_{\mathbb{R}^N} f U dx > 0.$$

By evaluating the functional in (2.2) on the sequence $u_n = nU$, we get

$$\frac{1}{2} \int_{\Omega} |\nabla u_n|^2 dx + \frac{1}{2} \int_{\Omega} V u_n^2 dx - \langle f, u_n \rangle = -n \langle f, U \rangle,$$

thus the functional is unbounded from below.

Lemma 2.4. Let $V_1, V_2 \in \mathcal{V}$ be two admissible potentials and let $u_1, u_2 \in W_0^{1,2}(\Omega)$ be solutions of (2.2). Then

$$(2.6) \quad \left| \int_{\Omega} V_1 u_1^2 dx - \int_{\Omega} V_2 u_2^2 dx \right| \leq C \|f\|_{W^{-1,2}(\Omega)} \|u_1 - u_2\|_{W_0^{1,2}(\Omega)},$$

where

$$C = 1 + \frac{T_N}{T_N - \|(V_1)_-\|_{L^{N/2}(\Omega)}} + \frac{T_N}{T_N - \|(V_2)_-\|_{L^{N/2}(\Omega)}}.$$

In particular $C = 3$ if V_1, V_2 are nonnegative.

Proof. From the respective PDEs, we obtain

$$\int_{\Omega} |\nabla u_i|^2 dx + \int_{\Omega} V_i u_i^2 dx = \langle f, u_i \rangle, \quad i = 1, 2,$$

so that

$$\left| \int_{\Omega} V_1 u_1^2 dx - \int_{\Omega} V_2 u_2^2 dx \right| \leq \left| \int_{\Omega} |\nabla u_1|^2 dx - \int_{\Omega} |\nabla u_2|^2 dx \right| + \|f\|_{W^{-1,2}(\Omega)} \|u_1 - u_2\|_{W_0^{1,2}(\Omega)}.$$

Finally, we observe that

$$\begin{aligned} \left| \|\nabla u_1\|_{L^2(\Omega)}^2 - \|\nabla u_2\|_{L^2(\Omega)}^2 \right| &= \left(\|\nabla u_1\|_{L^2(\Omega)} + \|\nabla u_2\|_{L^2(\Omega)} \right) \left| \|\nabla u_1\|_{L^2(\Omega)} - \|\nabla u_2\|_{L^2(\Omega)} \right| \\ &\leq \left(\|\nabla u_1\|_{L^2(\Omega)} + \|\nabla u_2\|_{L^2(\Omega)} \right) \|u_1 - u_2\|_{W_0^{1,2}(\Omega)}, \end{aligned}$$

then we can conclude by using (2.3). \square

Lemma 2.5. *Let $V_1, V_2 \in \mathcal{V}$ be two admissible potentials. Let $u_1, u_2 \in W_0^{1,2}(\Omega)$ be solutions of (2.2) such that*

$$u_i \in L^2(\Omega; V_1) \cap L^2(\Omega; V_2), \quad i = 1, 2.$$

Then we have

$$(2.7) \quad \mathcal{E}_f(V_1) - \mathcal{E}_f(V_2) = \frac{1}{2} \int_{\Omega} (V_1 - V_2) u_1 u_2 dx.$$

In particular there holds

$$(2.8) \quad |\mathcal{E}_f(V_1) - \mathcal{E}_f(V_2)| \leq \frac{1}{2} \left(\int_{\Omega} |V_1 - V_2| u_1^2 dx \right)^{1/2} \left(\int_{\Omega} |V_1 - V_2| u_2^2 dx \right)^{1/2}.$$

Proof. We first observe that, from the hypothesis on the potentials, we can use $u_1 - u_2$ as a test function for the equations solved by u_1 and u_2 , i.e.

$$(2.9) \quad \int_{\Omega} \langle \nabla u_i, \nabla \varphi \rangle dx + \int_{\Omega} V_i u_i \varphi dx = \langle f, \varphi \rangle, \quad \text{for every } \varphi \in W_0^{1,2}(\Omega) \cap L^2(\Omega; V_i), \quad i = 1, 2.$$

We have

$$\begin{aligned} \mathcal{E}_f(V_1) - \mathcal{E}_f(V_2) &= \frac{1}{2} \int_{\Omega} |\nabla u_1|^2 dx + \frac{1}{2} \int_{\Omega} V_1 u_1^2 dx - \langle f, u_1 \rangle \\ &\quad - \frac{1}{2} \int_{\Omega} |\nabla u_2|^2 dx - \frac{1}{2} \int_{\Omega} V_2 u_2^2 dx + \langle f, u_2 \rangle. \end{aligned}$$

On the other hand

$$\begin{aligned} \frac{1}{2} \int_{\Omega} |\nabla u_1|^2 dx - \frac{1}{2} \int_{\Omega} |\nabla u_2|^2 dx &= \frac{1}{2} \int_{\Omega} \langle \nabla u_1, \nabla (u_1 - u_2) \rangle dx \\ &\quad - \frac{1}{2} \int_{\Omega} \langle \nabla u_2, \nabla (u_2 - u_1) \rangle dx, \end{aligned}$$

thus by appealing to (2.9), we get

$$\begin{aligned} \frac{1}{2} \int_{\Omega} |\nabla u_1|^2 dx - \frac{1}{2} \int_{\Omega} |\nabla u_2|^2 dx &= -\frac{1}{2} \int_{\Omega} V_1 u_1 (u_1 - u_2) dx + \frac{1}{2} \langle f, u_1 - u_2 \rangle \\ &\quad + \frac{1}{2} \int_{\Omega} V_2 u_2 (u_2 - u_1) dx - \frac{1}{2} \langle f, u_2 - u_1 \rangle \\ &= -\frac{1}{2} \int_{\Omega} V_1 u_1^2 dx + \frac{1}{2} \int_{\Omega} V_2 u_2^2 dx + \langle f, u_1 - u_2 \rangle \\ &\quad + \frac{1}{2} \int_{\Omega} (V_1 - V_2) u_1 u_2 dx. \end{aligned}$$

This concludes the proof of (2.7). The estimate (2.8) just follows by applying Hölder inequality. \square

Remark 2.6. Observe that if $V_1, V_2 \in L^p(\Omega)$ with $p > 1$ and $u_1, u_2 \in L^{2p/(p-1)}(\Omega)$, then the hypotheses of the previous Lemma are verified and (2.8) gives the Lipschitz estimate

$$|\mathcal{E}_f(V_1) - \mathcal{E}_f(V_2)| \leq \frac{1}{2} \|V_1 - V_2\|_{L^p(\Omega)} \sqrt{\|u_1\|_{L^{2p/(p-1)}(\Omega)}} \sqrt{\|u_2\|_{L^{2p/(p-1)}(\Omega)}}.$$

By Proposition 2.1, the condition $u_1, u_2 \in L^{2p/(p-1)}(\Omega)$ is verified for example if

- $|\Omega| = +\infty$ and $p = N/2$;

- $\Omega = \omega \times \mathbb{R}$ is a waveguide and $p \geq N/2$;
- $|\Omega| < +\infty$ and $p \geq N/2$.

3. MAXIMIZATION PROBLEMS

In this section we fix $p > 1$ and we consider the optimization problem for potentials

$$(3.1) \quad \max_{V \in \mathcal{V}} \left\{ \mathcal{E}_f(V) : \int_{\Omega} |V|^p dx \leq 1 \right\}.$$

We also introduce the strictly convex functional

$$(3.2) \quad G_{p,f}(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2} \left(\int_{\Omega} |u|^{2p/(p-1)} dx \right)^{(p-1)/p} - \langle f, u \rangle, \quad u \in W_0^{1,2}(\Omega),$$

where it is intended that $G_{p,f}(u) = +\infty$ if $u \notin L^{2p/(p-1)}(\Omega)$. We recall the following existence result from [10]. We give the proof for the reader's convenience.

Proposition 3.1. *The problem (3.1) admits a solution and is equivalent to*

$$(3.3) \quad \max_{V \in \mathcal{V}} \left\{ \mathcal{E}_f(V) : V \geq 0, \int_{\Omega} V^p dx = 1 \right\}.$$

The solution V_0 is unique and is of the form

$$(3.4) \quad V_0 = \left(\int_{\Omega} |v_0|^{2p/(p-1)} dx \right)^{-1/p} v_0^{2/(p-1)},$$

where $v_0 \in W_0^{1,2}(\Omega) \cap L^{2p/(p-1)}(\Omega)$ is the unique minimizer of $G_{p,f}$. Moreover, we have

$$(3.5) \quad \mathcal{E}_f(V_0) = G_{p,f}(v_0).$$

Proof. We start by proving that we can restrict the optimization to positive potentials that saturate the constraint on the L^p norm. We have

$$\sup_{V \in \mathcal{V}} \left\{ \mathcal{E}_f(V) : V \geq 0, \int_{\Omega} |V|^p dx = 1 \right\} \leq \sup_{V \in \mathcal{V}} \left\{ \mathcal{E}_f(V) : \int_{\Omega} |V|^p dx \leq 1 \right\}.$$

On the other hand it is immediate to see that

$$\mathcal{E}_f(V) \leq \mathcal{E}_f \left(\frac{|V|}{\|V\|_{L^p(\Omega)}} \right), \quad \text{for every } V \in L^p(\Omega) \setminus \{0\} \text{ with } \int_{\Omega} |V|^p dx \leq 1,$$

thus the two suprema coincide.

In order to characterize the optimal potential V_0 , we observe that for every $u \in L^{2p/(p-1)}(\Omega)$ and every admissible potential, we get

$$(3.6) \quad \int_{\Omega} V u^2 dx \leq \left(\int_{\Omega} |u|^{2p/(p-1)} dx \right)^{(p-1)/p},$$

thanks to Hölder inequality. By appealing to the definition of the energy $\mathcal{E}_f(V)$, we then get

$$\mathcal{E}_f(V) \leq \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2} \left(\int_{\Omega} |u|^{2p/(p-1)} dx \right)^{(p-1)/p} - \langle f, u \rangle, \quad \forall u \in W_0^{1,2}(\Omega).$$

By taking the infimum on u , we obtain

$$\mathcal{E}_f(V) \leq \min_{u \in W_0^{1,2}(\Omega)} G_{p,f}(u), \quad \text{for every } V \text{ admissible.}$$

On the other hand, we see that if v_0 is a minimizer of $G_{p,f}$ and (V_0, v_0) achieves equality in (3.6), we have equality in the last inequality. By appealing to the equality cases in Hölder inequality, we get the characterization (3.4). \square

Remark 3.2. For the sake of completeness we observe that by a standard homogeneity argument

$$\mathcal{E}_f(V) = -\frac{1}{2} \sup_{u \in W_0^{1,2}(\Omega) \setminus \{0\}} \frac{\langle f, u \rangle^2}{\int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} V u^2 dx}.$$

By using (3.6) we can infer

$$\mathcal{E}_f(V) \leq -\frac{1}{2} \sup_{u \in W_0^{1,2}(\Omega) \setminus \{0\}} \frac{\langle f, u \rangle^2}{\int_{\Omega} |\nabla u|^2 dx + \left(\int_{\Omega} |u|^{2p/(p-1)} dx \right)^{(p-1)/p}},$$

for every $V \in L^p(\Omega)$ with unit norm. This implies that

$$\sup \left\{ \mathcal{E}_f(V) : \int_{\Omega} |V|^p dx = 1 \right\} = -\frac{1}{2} \sup_{u \in W_0^{1,2}(\Omega) \setminus \{0\}} \frac{\langle f, u \rangle^2}{\int_{\Omega} |\nabla u|^2 dx + \left(\int_{\Omega} |u|^{2p/(p-1)} dx \right)^{(p-1)/p}},$$

so that $\mathcal{E}_f(V_0)$ is related to the best constant in a Poincaré-Sobolev type inequality, i.e.

$$\int_{\Omega} |\nabla u|^2 dx + \left(\int_{\Omega} |u|^{2p/(p-1)} dx \right)^{(p-1)/p} \geq \frac{1}{2|\mathcal{E}_f(V_0)|} \langle u, f \rangle^2, \quad \forall u \in W_0^{1,2}(\Omega),$$

with equality holding if and only if u is proportional to v_0 .

4. STABILITY FOR MAXIMIZATION PROBLEMS

In what follows c_1 will denote the constant

$$c_1 := \left(\int_{\Omega} |v_0|^{2p/(p-1)} dx \right)^{(p-1)/2p},$$

where v_0 is the unique minimizer of $G_{p,f}$. In this section we prove a quantitative improvement of the inequality

$$\mathcal{E}_f(V_0) \geq \mathcal{E}_f(V), \quad \text{for every } V \in \mathcal{V} \text{ such that } \int_{\Omega} |V|^p dx \leq 1.$$

At this aim, we need the following result, see [12, Theorem 3.1] for a proof. An earlier related result could be found in [14, Proposition 2.6].

Lemma 4.1 (Quantitative Hölder inequality). *Let $2 \leq q < \infty$ and $q' = q/(q-1)$. For every $f \in L^q(\Omega)$ and $g \in L^{q'}(\Omega)$ such that $\|f\|_{L^q(\Omega)} = \|g\|_{L^{q'}(\Omega)} = 1$, we have*

$$(4.1) \quad \left| \int_{\Omega} f g dx \right| \leq 1 - \frac{q' - 1}{4} \left\| |f|^{q-2} f - g \right\|_{L^{q'}(\Omega)}^2,$$

and

$$(4.2) \quad \left| \int_{\Omega} f g dx \right| \leq 1 - \frac{1}{q 2^{q-1}} \left\| f - |g|^{q'-2} g \right\|_{L^q(\Omega)}^q.$$

Remark 4.2. In the case $q = \infty$ no quantitative inequality of the previous kind may hold. In fact, by taking $\Omega = (0, 1)$ and the functions

$$f_n = 1_{[0, 1-1/n]} \quad \text{and} \quad g = \frac{1}{2\sqrt{x}},$$

we obtain

$$\lim_{n \rightarrow \infty} \left(\left| \int_{\Omega} f_n g dx \right| - 1 \right) = 0 \quad \text{while} \quad \left\| f_n - \frac{g}{|g|} \right\|_{L^\infty(\Omega)} = \|f_n - 1\|_{L^\infty(\Omega)} = 1$$

4.1. Stability of the potentials. This is the main result of this section.

Theorem 4.3 (Stability of maximal potentials). *Let V_0 be the optimal potential achieving the maximum in (3.1). Then for every $V \in \mathcal{V}$ such that $\|V\|_{L^p(\Omega)} \leq 1$ we have*

$$(4.3) \quad \mathcal{E}_f(V_0) - \mathcal{E}_f(V) \geq \sigma'_M \left\| |V|^{p-2} V - V_0^{p-1} \right\|_{L^{p'}(\Omega)}^2, \quad \text{for } p \geq 2,$$

and

$$(4.4) \quad \mathcal{E}_f(V_0) - \mathcal{E}_f(V) \geq \sigma''_M \|V - V_0\|_{L^p(\Omega)}^2, \quad \text{for } 1 < p < 2,$$

where $\sigma'_M > 0$ and $\sigma''_M > 0$ are two constants depending only on p and c_1 (see Remark 4.4 below).

Proof. We start observing that by hypothesis

$$\|V - V_0\|_{L^p(\Omega)} \leq 2, \quad \text{and} \quad \left\| |V|^{p-2} V - |V_0|^{p-2} V_0 \right\|_{L^{p'}(\Omega)} \leq 2,$$

thus we can always suppose

$$(4.5) \quad \mathcal{E}_f(V_0) - \mathcal{E}_f(V) \leq \min \left\{ \frac{c_1^2}{4}, 1 \right\},$$

because otherwise (4.3) and (4.4) are trivially true, with constants

$$\sigma'_M = \sigma''_M = \frac{1}{4} \min \left\{ \frac{c_1^2}{4}, 1 \right\}.$$

By using v_0 as a test function in the variational problem defining $\mathcal{E}_f(V)$ and recalling the definition (3.2) of $G_{p,f}$, we get

$$\begin{aligned} \mathcal{E}_f(V) &\leq G_{p,f}(v_0) + \frac{1}{2} \left[\int_{\Omega} V v_0^2 dx - \left(\int_{\Omega} |v_0|^{2p/(p-1)} dx \right)^{(p-1)/p} \right] \\ &= \mathcal{E}_f(V_0) + \frac{1}{2} \left[\int_{\Omega} V v_0^2 dx - \left(\int_{\Omega} |v_0|^{2p/(p-1)} dx \right)^{(p-1)/p} \right]. \end{aligned}$$

The optimal potential V_0 and v_0 are linked through (3.4), thus by substituting above we get

$$\mathcal{E}_f(V) - \mathcal{E}_f(V_0) \leq \frac{c_1^2}{2} \int_{\Omega} (V - V_0) V_0^{p-1} dx = \frac{c_1^2}{2} \left[\int_{\Omega} V V_0^{p-1} dx - 1 \right].$$

From the previous inequality we obtain

$$(4.6) \quad \mathcal{E}_f(V_0) - \mathcal{E}_f(V) \geq \frac{c_1^2}{2} \left[1 - \left(\int_{\Omega} |V|^p dx \right)^{1/p} \right] + \frac{c_1^2}{2} \left[\left(\int_{\Omega} |V|^p dx \right)^{1/p} - \int_{\Omega} V V_0^{p-1} dx \right].$$

The two terms inside the square brackets are both positive, since $\|V\|_{L^p(\Omega)} \leq \|V_0^{p-1}\|_{L^{p'}(\Omega)} = 1$. The previous estimate in particular implies that

$$(4.7) \quad \left(\int_{\Omega} |V|^p dx \right)^{1/p} \geq \frac{1}{2},$$

since otherwise we would contradict (4.5). We now distinguish two cases.

Case $p \geq 2$. By applying (4.1) with the choices

$$f = \frac{V}{\|V\|_{L^p(\Omega)}}, \quad g = V_0^{p-1}, \quad q = p \quad \text{and} \quad q' = p',$$

we get

$$(4.8) \quad \left(\int_{\Omega} |V|^p dx \right)^{1/p} - \int_{\Omega} V V_0^{p-1} dx \geq \frac{p' - 1}{8} \left\| \frac{|V|^{p-2} V}{\|V\|_{L^p(\Omega)}^{p-1}} - V_0^{p-1} \right\|_{L^{p'}(\Omega)}^2,$$

where we used (4.7) to estimate the L^p norm of V from below. We now observe that by the triangle inequality and convexity of $t \mapsto t^2$, we get

$$\begin{aligned} \left\| \frac{|V|^{p-2} V}{\|V\|_{L^p(\Omega)}^{p-1}} - V_0^{p-1} \right\|_{L^{p'}(\Omega)}^2 &\geq \frac{1}{2} \left\| |V|^{p-2} V - V_0^{p-1} \right\|_{L^{p'}(\Omega)}^2 - \left\| \frac{|V|^{p-2} V}{\|V\|_{L^p(\Omega)}^{p-1}} - |V|^{p-2} V \right\|_{L^{p'}(\Omega)}^2 \\ &= \frac{1}{2} \left\| |V|^{p-2} V - V_0^{p-1} \right\|_{L^{p'}(\Omega)}^2 - \left| 1 - \|V\|_{L^p(\Omega)}^{p-1} \right|^2 \\ &\geq \frac{1}{2} \left\| |V|^{p-2} V - V_0^{p-1} \right\|_{L^{p'}(\Omega)}^2 - (p-1)^2 \left(\frac{2}{c_1^2} \right)^2 \left(\mathcal{E}_f(V_0) - \mathcal{E}_f(V) \right)^2, \end{aligned}$$

where we used that for $p \geq 2$

$$1 - t^{p-1} \leq (p-1)(1-t), \quad \text{for every } 0 \leq t \leq 1,$$

and (4.6) in the last inequality. By inserting this information in (4.8), combining with (4.6) and using that $\mathcal{E}_f(V_0) - \mathcal{E}_f(V) \leq 1$, we end up with (4.3).

Case $1 < p < 2$. By applying (4.1) this time with the choices

$$f = V_0^{p-1}, \quad g = \frac{V}{\|V\|_{L^p(\Omega)}}, \quad q = p' \quad \text{and} \quad q' = p,$$

we get

$$(4.9) \quad \left(\int_{\Omega} |V|^p dx \right)^{1/p} - \int_{\Omega} V V_0^{p-1} dx \geq \frac{p-1}{8} \left\| \frac{V}{\|V\|_{L^p(\Omega)}} - V_0 \right\|_{L^p(\Omega)}^2,$$

where we used again (4.7). We can estimate the remainder term as before

$$\begin{aligned} \left\| \frac{V}{\|V\|_{L^p(\Omega)}} - V_0 \right\|_{L^p(\Omega)}^2 &\geq \frac{1}{2} \|V - V_0\|_{L^p(\Omega)}^2 - \left\| \frac{V}{\|V\|_{L^p(\Omega)}} - V \right\|_{L^p(\Omega)}^2 \\ &= \frac{1}{2} \|V - V_0\|_{L^p(\Omega)}^2 - |1 - \|V\|_{L^p(\Omega)}|^2 \\ &\geq \frac{1}{2} \|V - V_0\|_{L^p(\Omega)}^2 - \left(\frac{2}{c_1^2} \right)^2 \left(\mathcal{E}_f(V_0) - \mathcal{E}_f(V) \right)^2, \end{aligned}$$

where we used again (4.6) in the last inequality. We can obtain the desired result by combining (4.6), (4.9) and the previous estimate. \square

Remark 4.4. From the proof, we can see that a possible value for σ'_M is

$$\sigma'_M = \frac{1}{4} \min \left\{ (p'-1) \frac{c_1^4}{8c_1^2 + 2(p-1)}, 1, \frac{c_1^2}{4} \right\}, \quad p \geq 2,$$

while for σ''_M we could take

$$\sigma''_M = \frac{1}{4} \min \left\{ (p-1) \frac{c_1^4}{8c_1^2 + 2(p-1)}, 1, \frac{c_1^2}{4} \right\}, \quad 1 < p < 2.$$

Remark 4.5. We point out that we could have used (4.2) in place of (4.1). In this way, one could obtain stability estimates of the type

$$(4.10) \quad \mathcal{E}_f(V_0) - \mathcal{E}_f(V) \geq \tilde{\sigma}'_M \|V - V_0\|_{L^p(\Omega)}^p, \quad \text{for } p \geq 2,$$

and

$$\mathcal{E}_f(V_0) - \mathcal{E}_f(V) \geq \tilde{\sigma}''_M \left\| |V|^{p-2} V - V_0^{p-1} \right\|_{L^{p'}(\Omega)}^{p'}, \quad \text{for } 1 < p < 2.$$

We also notice that *these estimates are asymptotically worse* than (4.3) and (4.4). Indeed, when V is of the form $V_\varepsilon = V_0 + \varepsilon \psi$, for $\varepsilon \ll 1$ and $\psi \in L^p(\Omega)$, it is not difficult to see that

$$\|V_\varepsilon - V_0\|_{L^p(\Omega)} \simeq \varepsilon \quad \text{and} \quad \left\| |V_\varepsilon|^{p-2} V_\varepsilon - V_0^{p-1} \right\|_{L^{p'}(\Omega)} \simeq \varepsilon,$$

so that we have

$$\| |V_\varepsilon|^{p-2} V_\varepsilon - V_0^{p-1} \|_{L^{p'}(\Omega)}^2 \gg \|V_\varepsilon - V_0\|_{L^p(\Omega)}^p \quad \text{if } p \geq 2,$$

$$\|V_\varepsilon - V_0\|_{L^p(\Omega)}^2 \gg \| |V_\varepsilon|^{p-2} V_\varepsilon - V_0^{p-1} \|_{L^{p'}(\Omega)}^{p'} \quad \text{if } 1 < p < 2.$$

4.2. Stability of the state functions. We have the following stability result for the minimization of $G_{p,f}$.

Proposition 4.6. *Let $1 < p < \infty$ and let v_0 be the unique minimizer of $G_{p,f}$ defined by (3.2), then for every $u \in W_0^{1,2}(\Omega)$ we have*

$$(4.11) \quad G_{p,f}(u) - G_{p,f}(v_0) \geq \frac{1}{2} \|u - v_0\|_{W_0^{1,2}(\Omega)}^2.$$

Proof. We first observe that if $u \notin L^{2p/(p-1)}(\Omega)$, then $G_{p,f}(u) = +\infty$ and (4.11) trivially holds. Thus, let us take $u \in W_0^{1,2}(\Omega) \cap L^{2p/(p-1)}(\Omega)$ with $u \neq v_0$ and set

$$d = \|u - v_0\|_{W_0^{1,2}(\Omega)} \quad \text{and} \quad \varphi = \frac{u - v_0}{d}.$$

Then u can be written as $u = v_0 + d\varphi$ and φ has unitary norm in $W_0^{1,2}(\Omega)$. We then get

$$\begin{aligned} G_{p,f}(u) - G_{p,f}(v_0) &\geq d \left[\int_{\Omega} \langle \nabla v_0, \nabla \varphi \rangle dx + \left(\int_{\Omega} |v_0|^{2p/(p-1)} dx \right)^{-1/p} \int_{\Omega} |v_0|^{2/(p-1)} v_0 \varphi dx - \langle f, \varphi \rangle \right] \\ &\quad + \frac{d^2}{2} \int_{\Omega} |\nabla \varphi|^2 dx, \end{aligned}$$

where we used the convexity of the C^1 map

$$\Phi(s) = \|v_0 + s\varphi\|_{L^{2p/(p-1)}(\Omega)}^2 = \left(\int_{\Omega} |v_0 + s\varphi|^{2p/(p-1)} dx \right)^{(p-1)/p}, \quad s \in \mathbb{R},$$

so that

$$\Phi(d) \geq \Phi(0) + \Phi'(0)d.$$

It is now sufficient to observe that

$$\int_{\Omega} \langle \nabla v_0, \nabla \varphi \rangle dx + \left(\int_{\Omega} |v_0|^{2p/(p-1)} dx \right)^{-1/p} \int_{\Omega} |v_0|^{2/(p-1)} v_0 \varphi dx - \langle f, \varphi \rangle = 0,$$

by minimality of v_0 , thus we directly get (6.2). \square

As a consequence of the previous result, we get that if (u, V) is almost realizing the equality in (3.6), then u is near to the optimizer v_0 in the $W_0^{1,2}$ norm.

Corollary 4.7. *Let V be an admissible potential for (3.1) and u a corresponding energy function. Then*

$$(4.12) \quad \|u - v_0\|_{W_0^{1,2}(\Omega)}^2 \leq \left[\left(\int_{\Omega} |u|^{2p/(p-1)} dx \right)^{(p-1)/p} - \int_{\Omega} V u^2 dx \right].$$

Proof. We already observed that

$$\mathcal{E}_f(V) \leq G_{p,f}(v), \quad \text{for every } v \in W_0^{1,2}(\Omega).$$

By taking the energy function v_0 corresponding to V_0 , we get

$$\mathcal{E}_f(V) \leq G_{p,f}(v_0) \leq G_{p,f}(u) - \frac{1}{2} \|u - v_0\|_{W_0^{1,2}(\Omega)}^2,$$

where we used (4.11). Thus we have

$$\frac{1}{2} \|u - v_0\|_{W_0^{1,2}(\Omega)}^2 \leq G_{p,f}(u) - \mathcal{E}_f(V) = \frac{1}{2} \left[\left(\int_{\Omega} |u|^{2p/(p-1)} dx \right)^{(p-1)/p} - \int_{\Omega} V u^2 dx \right],$$

which concludes the proof. \square

In general, for an admissible potential V the corresponding energy function is not in $L^{2p/(p-1)}(\Omega)$. When this is the case, we can infer stability of the energy functions as well.

Proposition 4.8. *Let V be admissible in (3.1). If a corresponding energy function u belongs to $L^{2p/(p-1)}(\Omega)$, then*

$$(4.13) \quad \left(\mathcal{E}_f(V_0) - \mathcal{E}_f(V) \right)^{\frac{1}{\vartheta(p)}} \|u - v_0\|_{L^{2p/(p-1)}(\Omega)} \geq c \|u - v_0\|_{W_0^{1,2}(\Omega)}^2,$$

where $\vartheta(p) = \max\{2, p\}$ and $c > 0$ is a constant depending only on c_1 , p and $\|V_-\|_{L^{N/2}(\Omega)}$.

Proof. We first observe that since $u \in L^{2p/(p-1)}(\Omega)$ and $u_0 \in L^{2p/(p-1)}(\Omega)$ as well, we have

$$\psi := u - v_0 \in W_0^{1,2}(\Omega) \cap L^2(\Omega; V).$$

The function ψ verifies

$$\int_{\Omega} \langle \nabla \psi, \nabla \varphi \rangle dx + \int_{\Omega} (V - V_0) v_0 \varphi dx + \int_{\Omega} V \psi \varphi dx = 0,$$

for every $\varphi \in W_0^{1,2}(\Omega)$. By using ψ itself as a test function, we get

$$\begin{aligned} \int_{\Omega} |\nabla \psi|^2 dx + \int_{\Omega} V \psi^2 dx &\leq \int_{\Omega} |V - V_0| |v_0| |\psi| dx \\ &\leq \left(\int_{\Omega} |V - V_0|^p dx \right)^{1/p} \left(\int_{\Omega} (|v_0| |\psi|)^{p/(p-1)} dx \right)^{(p-1)/p} \\ &\leq \|V - V_0\|_{L^p(\Omega)} \|v_0\|_{L^{2p/(p-1)}(\Omega)} \|\psi\|_{L^{2p/(p-1)}(\Omega)}. \end{aligned}$$

By noticing that from (2.4)

$$\int_{\Omega} |\nabla \psi|^2 dx + \int_{\Omega} V \psi^2 dx \geq \left(1 - \frac{\|V_-\|_{L^{N/2}(\Omega)}}{T_N} \right) \int_{\Omega} |\nabla \psi|^2 dx,$$

and recalling that $\psi = u - v_0$, we get

$$\left(1 - \frac{\|V_-\|_{L^{N/2}(\Omega)}}{T_N} \right) \|u - v_0\|_{W_0^{1,2}(\Omega)}^2 \leq c_1 \|V - V_0\|_{L^p(\Omega)} \|u - v_0\|_{L^{2p/(p-1)}(\Omega)}.$$

Appealing to (4.4) (for $1 < p < 2$) or to (4.10) (for $p \geq 2$), we then get the conclusion. \square

Remark 4.9. Observe that for $p = N/2$, we have $2p/(p-1) = 2^*$ and (4.13) simply becomes

$$\left(\mathcal{E}_f(V_0) - \mathcal{E}_f(V) \right)^{\frac{1}{\vartheta(p)}} \geq c \|u - v_0\|_{W_0^{1,2}(\Omega)},$$

by Sobolev inequality, possibly with a different constant $c > 0$. When Ω has a finite measure or is a waveguide and $p \geq N/2$, by Proposition 2.1 we can always assure that the energy function u belongs to $L^{2p/(p-1)}(\Omega)$ and thus we have a similar stability estimate in these cases as well.

5. MINIMIZATION PROBLEMS

In this section we consider, for a fixed $p > 0$, the minimization problem

$$(5.1) \quad \inf \left\{ \mathcal{E}_f(V) : V \geq 0, \int_{\Omega} \frac{1}{V^p} dx \leq 1 \right\}.$$

Remark 5.1. Observe that this time, it is not clear whether the minimization problem on \mathcal{V} without sign hypothesis, i.e.

$$\inf_{V \in \mathcal{V}} \left\{ \mathcal{E}_f(V) : \int_{\Omega} \frac{1}{|V|^p} dx \leq 1 \right\}.$$

is well-posed or not, since it may happen that no admissible $V \leq 0$ exist (for example if Ω is unbounded). If an optimal potential W exists for the previous problem, this should be such that $W_+ \not\equiv 0$ and $W_- \not\equiv 0$. This seems to be an interesting issue, which we leave for future research.

We collect a couple of technical results which are needed in the sequel.

Lemma 5.2. *Let $V \geq 0$ be such that $V^{-1} \in L^p(\Omega)$. Then we have*

$$(5.2) \quad \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \leq \left(\int_{\Omega} V u^2 dx \right) \left(\int_{\Omega} \frac{1}{V^p} dx \right)^{1/p}, \quad \text{for every } u \in C_0^\infty(\Omega).$$

In particular we have the continuous embedding $L^2(\Omega; V) \subset L^{2p/(p+1)}(\Omega)$.

Proof. If $u \notin L^2(\Omega; V)$, there is nothing to prove. So let us assume that the first integral in the right-hand side of (5.2) is finite. By Hölder inequality with exponents $q = (p+1)/p$ and $q' = p+1$ we have

$$\int_{\Omega} |u|^{2p/(p+1)} dx = \int_{\Omega} |u|^{2p/(p+1)} \frac{V^{p/(p+1)}}{V^{p/(p+1)}} dx \leq \left(\int_{\Omega} V u^2 dx \right)^{p/(p+1)} \left(\int_{\Omega} \frac{1}{V^p} dx \right)^{1/(p+1)},$$

which concludes the proof. \square

Remark 5.3. By standard interpolation in Lebesgue spaces, under the hypothesis of the previous result we have the continuous embedding

$$L^2(\Omega; V) \subset L^s(\Omega),$$

for every $2p/(p+1) \leq s \leq 2^*$. In particular $L^2(\Omega; V)$ is embedded into $L^2(\Omega)$, since $2p/(p+1)$ is always strictly less than 2. It is then not difficult to show that the operator $-\Delta + V$ has a discrete spectrum on $L^2(\Omega)$.

The following energy estimate will be needed in the sequel.

Lemma 5.4. *Let $f \in W^{-1,2}(\Omega)$, V be an admissible potential for (5.1) and u its energy function. Then we have*

$$(5.3) \quad \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \leq \frac{1}{2} \|f\|_{W^{-1,2}(\Omega)}^2.$$

Proof. By using the equation and Young inequality, we have

$$\int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} V u^2 dx = \langle f, u \rangle \leq \frac{1}{2} \|f\|_{W^{-1,2}(\Omega)}^2 + \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx.$$

From the previous we obtain

$$\int_{\Omega} V u^2 dx \leq \frac{1}{2} \|f\|_{W^{-1,2}(\Omega)}^2,$$

then it is sufficient to apply (5.2). \square

Let $1 < p < \infty$, in what follows we set for simplicity

$$J_{p,f}(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2} \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} - \langle f, u \rangle \quad \forall u \in W_0^{1,2}(\Omega),$$

where it is intended that $J_{p,f}(u) = +\infty$ if $u \notin L^{2p/(p+1)}(\Omega)$. Again, it is not difficult to see that $J_{p,f}$ admits a unique minimizer. We recall the following result from [10].

Proposition 5.5. *The problem (5.1) admits a unique solution U_0 of the form*

$$(5.4) \quad U_0 = \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{1/p} |u_0|^{-2/(p+1)},$$

where $u_0 \in W_0^{1,2}(\Omega) \cap L^{2p/(p+1)}(\Omega)$ is the unique minimizer of $J_{p,f}$.

Proof. Let u be the energy function corresponding to V , then we have

$$\mathcal{E}_f(V) = J_{p,f}(u) + \frac{1}{2} \left[\int_{\Omega} V u^2 - \left(\int_{\Omega} |u|^{2p/(p+1)} \right)^{(p+1)/p} \right].$$

By using (5.2) and the minimality of u_0 , we get

$$\mathcal{E}_f(V) \geq J_{p,f}(u_0) = \min_{u \in W_0^{1,2}(\Omega)} J_{p,f}(u).$$

By appealing again to the equality cases in Hölder inequality, we get the characterization of the optimal potential V_0 . \square

Remark 5.6. As in Remark 3.2, by applying (5.2) we can infer

$$\mathcal{E}_f(V) \geq -\frac{1}{2} \sup_{u \in W_0^{1,2}(\Omega) \setminus \{0\}} \frac{\langle f, u \rangle^2}{\int_{\Omega} |\nabla u|^2 dx + \left(\int_{\Omega} |u|^{2p/(p+1)} \right)^{(p+1)/p}},$$

for every V admissible, so that

$$\inf \left\{ \mathcal{E}_f(V) : V \geq 0, \int_{\Omega} V^{-p} dx \leq 1 \right\} = -\frac{1}{2} \sup_{u \in W_0^{1,2}(\Omega) \setminus \{0\}} \frac{\langle f, u \rangle^2}{\int_{\Omega} |\nabla u|^2 dx + \left(\int_{\Omega} |u|^{2p/(p+1)} \right)^{(p+1)/p}}.$$

Thus we have

$$\int_{\Omega} |\nabla u|^2 dx + \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \geq \frac{1}{2|\mathcal{E}_f(U_0)|} \langle u, f \rangle^2, \quad u \in W_0^{1,2}(\Omega) \cap L^{2p/(p+1)}(\Omega),$$

with equality if and only if u is proportional to u_0 .

6. STABILITY FOR MINIMIZATION PROBLEMS

In this section we take $p > 1$ and we still denote by u_0 the unique minimizer of $J_{p,f}$. We also set

$$(6.1) \quad c_2 := \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{(p+1)/2p}.$$

6.1. Preliminary results. We start with a stability result for the minimization of $J_{p,f}$. The proof is the same as that of Lemma 4.6, thus we omit it.

Proposition 6.1. *Let $1 < p < \infty$ and let u_0 be the unique minimizer of $J_{p,f}$. Then for every $u \in W_0^{1,2}(\Omega)$ we have*

$$(6.2) \quad J_{p,f}(u) - J_{p,f}(u_0) \geq \frac{1}{2} \|u - u_0\|_{W_0^{1,2}(\Omega)}^2.$$

We also need the following result, asserting that the energy gap controls the difference of the $L^{2p/(p+1)}(\Omega)$ norms.

Lemma 6.2. *Let $f \in W^{-1,2}(\Omega)$ and let V be an admissible potential for (5.1). If we suppose that $\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \leq 1$, then*

$$(6.3) \quad \sqrt{\mathcal{E}_f(V) - \mathcal{E}_f(U_0)} \geq c_3 \left| \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{(p+1)/p} - \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \right|,$$

for a constant $c_3 > 0$ depending only on $\|f\|_{W^{-1,2}(\Omega)}$.

Proof. We first observe that

$$(6.4) \quad \mathcal{E}_f(V) - \mathcal{E}_f(U_0) = [J_{p,f}(u) - J_{p,f}(u_0)] + \frac{1}{2} \left[\int_{\Omega} V u^2 dx - \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \right],$$

and both terms inside the square brackets are positive. In particular we get

$$\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \geq \frac{1}{2} \left[\int_{\Omega} V u^2 dx - \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \right].$$

By using the estimate on the weighted L^2 norms (2.6), we then get

$$(6.5) \quad \mathcal{E}_f(V) - \mathcal{E}_f(U_0) \geq \frac{1}{2} \left[\int_{\Omega} U_0 u_0^2 dx - \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \right] - \frac{3}{2} \|f\|_{W^{-1,2}(\Omega)} \|u - u_0\|_{W^{1,2}(\Omega)}.$$

We now use that by (5.4) U_0 and u_0 are linked through

$$U_0 = \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{1/p} |u_0|^{-2/(p+1)},$$

so that

$$(6.6) \quad \int_{\Omega} U_0 u_0^2 dx = \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{(p+1)/p}.$$

If we use this in (6.5), we end up with

$$\begin{aligned} \mathcal{E}_f(V) - \mathcal{E}_f(U_0) &+ \frac{3}{2} \|f\|_{W^{-1,2}(\Omega)} \|u - u_0\|_{W^{1,2}(\Omega)} \\ &\geq \frac{1}{2} \left[\left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{(p+1)/p} - \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \right]. \end{aligned}$$

By recalling (6.2), (6.4) and using the hypothesis $\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \leq 1$, we then get

$$\sqrt{\mathcal{E}_f(V) - \mathcal{E}_f(U_0)} \geq \frac{1}{C} \left[\left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{(p+1)/p} - \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \right].$$

On the other hand, we have

$$\left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} - \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{(p+1)/p} \leq \int_{\Omega} V u^2 dx - \int_{\Omega} U_0 u_0^2 dx,$$

where we used (5.2) and (6.6). If we now apply (2.6), we get

$$\left[\left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} - \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{(p+1)/p} \right] \leq 3 \|f\|_{W^{-1,2}(\Omega)} \|u - u_0\|_{W_0^{1,2}(\Omega)},$$

and again we can conclude thanks to (6.4) and (6.2). \square

Remark 6.3. A closer inspection of the previous proof ensures that we can take

$$c_3 := \min \left\{ \frac{\sqrt{2}}{\sqrt{2} + 3 \|f\|_{W^{-1,2}(\Omega)}}, \frac{1}{3\sqrt{2} \|f\|_{W^{-1,2}(\Omega)}} \right\}.$$

The following result guarantees that it is sufficient to prove stability for potentials saturating the constraint $\int_{\Omega} V^{-p} \leq 1$.

Lemma 6.4 (Reduction Lemma). *Let V be admissible in (5.1) and such that $\int_{\Omega} V^{-p} dx < 1$. Let us suppose that its energy function u satisfies*

$$(6.7) \quad \int_{\Omega} V u^2 dx \geq \beta > 0.$$

Then there exists a potential $U \geq 0$ with $\int_{\Omega} U^{-p} dx = 1$ such that

$$(6.8) \quad \mathcal{E}_f(V) - \mathcal{E}_f(U_0) \geq \mathcal{E}_f(U) - \mathcal{E}_f(U_0),$$

and

$$(6.9) \quad \left\| \frac{1}{V} - \frac{1}{U_0} \right\|_{L^p(\Omega)} \leq \left\| \frac{1}{U} - \frac{1}{U_0} \right\|_{L^p(\Omega)} + \frac{2}{\beta} \left(\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \right).$$

Proof. Let $\lambda = \|V^{-1}\|_{L^p(\Omega)} < 1$, then we set $U = \lambda V$. It is clear that $U < V$, so that $\mathcal{E}_f(U) \leq \mathcal{E}_f(V)$ by the definition of the energy and the first property (6.8) follows. In order to prove the second, we observe that by (6.4), we have

$$\begin{aligned} \mathcal{E}_f(V) - \mathcal{E}_f(U_0) &\geq \frac{1}{2} \left[\int_{\Omega} V u^2 dx - \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \right] \\ &\geq \frac{1}{2} \int_{\Omega} V u^2 dx \left[1 - \left(\int_{\Omega} V^{-p} dx \right)^{1/p} \right] = \frac{1}{2} \left(\int_{\Omega} V u^2 dx \right) |1 - \lambda|, \end{aligned}$$

where we also used (5.2) in the second inequality. By using the hypothesis on u and the definition of U , we get

$$|1 - \lambda| \leq \frac{2}{\beta} \left(\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \right),$$

and since by the triangle inequality

$$\left\| \frac{1}{V} - \frac{1}{U_0} \right\|_{L^p(\Omega)} \leq \left\| \frac{1}{U} - \frac{1}{U_0} \right\|_{L^p(\Omega)} + |1 - \lambda|,$$

we get the desired conclusion. \square

Finally, the following very simple estimate will be quite useful.

Lemma 6.5. *Let $1 < r, s < \infty$ and $g, g_0 \in L^r(\Omega) \cap L^s(\Omega)$. Then we have*

$$(6.10) \quad \|g_0\|_{L^r(\Omega)} \leq \|g\|_{L^r(\Omega)} + \|g - g_0\|_{L^s(\Omega)} \frac{\| |g_0|^{r-1} \|_{L^{s'}(\Omega)}}{\|g_0\|_{L^r(\Omega)}^{r-1}}.$$

Proof. We can suppose that $|g_0|^{r-1} \in L^{s'}(\Omega)$, otherwise there is nothing to prove. For every $\varphi \in L^{r'}(\Omega) \cap L^{s'}(\Omega)$ we have

$$\begin{aligned} \left| \int_{\Omega} g_0 \varphi dx \right| &\leq \int_{\Omega} |g - g_0| |\varphi| dx + \left| \int_{\Omega} g \varphi dx \right| \\ &\leq \|g - g_0\|_{L^s(\Omega)} \|\varphi\|_{L^{s'}(\Omega)} + \|g\|_{L^r(\Omega)} \|\varphi\|_{L^{r'}(\Omega)}. \end{aligned}$$

If we now choose

$$\varphi = |g_0|^{r-2} g_0,$$

and then simplify by $\|g_0\|_{L^r(\Omega)}^{r-1}$ on both sides, we obtain (6.10). \square

6.2. Stability of the potentials. The following is the main result of this section, which is proved under the integrability assumption (6.11) on u_0 . For a discussion on this hypothesis, we refer the reader to Remark 6.7 below.

Theorem 6.6 (Stability of minimal potentials). *Let U_0 be the optimal potential achieving the minimum in (5.1). Let us suppose that the optimal function u_0 is such that*

$$(6.11) \quad c_4 := \left\| |u_0|^{(p-1)/(p+1)} \right\|_{L^{(2^*)}'(\Omega)} < +\infty.$$

Then for every positive potential V such that $\|1/V\|_{L^p(\Omega)} \leq 1$ we have

$$(6.12) \quad \mathcal{E}_f(V) - \mathcal{E}_f(U_0) \geq \sigma_m \left\| \frac{1}{V} - \frac{1}{U_0} \right\|_{L^p(\Omega)}^{2p(p+1)/(p-1)},$$

for a constant $\sigma_m > 0$ depending only on N, p, c_2, c_4 and $\|f\|_{W^{-1,2}(\Omega)}$ (see Remark 6.8 below).

Proof. We divide the proof into various steps.

Reduction step. Let V be a potential admissible in (5.1). We set

$$(6.13) \quad c_5 = \min \left\{ 1, \left(\frac{c_2^2 c_3}{2} \right)^2 \right\},$$

where we recall that c_2 is the $L^{2p/(p+1)}(\Omega)$ norm of u_0 and c_3 is the constant in (6.3). Since by hypothesis

$$\left\| \frac{1}{V} - \frac{1}{U_0} \right\|_{L^p(\Omega)} \leq 2,$$

we can always assume

$$\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \leq c_5,$$

otherwise (6.12) is trivially true with the constant $\sigma_m = c_5 4^{-p(p+1)/(p-1)}$. Under this assumption, by definition of c_5 and Lemma 6.2 we have that the energy function v of V verifies

$$\left(\int_{\Omega} |v|^{2p/(p+1)} dx \right)^{(p+1)/p} \geq \frac{1}{2} \left(\int_{\Omega} |u_0|^{2p/(p+1)} dx \right)^{(p+1)/p} = \frac{c_2^2}{2}.$$

This in turn implies that v verifies (6.7) with $\beta = c_2^2/2$ thanks to (5.2). From the Reduction Lemma 6.4 we thus obtain that there exists a positive potential U with $\|1/U\|_{L^p(\Omega)} = 1$ such that

$$\left\| \frac{1}{V} - \frac{1}{U_0} \right\|_{L^p(\Omega)} \leq \left\| \frac{1}{U} - \frac{1}{U_0} \right\|_{L^p(\Omega)} + \frac{4}{c_2^2} (\mathcal{E}_f(V) - \mathcal{E}_f(U_0))$$

and

$$\mathcal{E}_f(U) - \mathcal{E}_f(U_0) \leq \mathcal{E}_f(V) - \mathcal{E}_f(U_0).$$

We are going to prove the stability estimate (6.12) for the potential U . Observe that since the energy gap has decreased, we still have

$$(6.14) \quad \mathcal{E}_f(U) - \mathcal{E}_f(U_0) \leq c_5,$$

and thus again

$$(6.15) \quad \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \geq \frac{c_2^2}{2}.$$

where u is now the energy function of U . From (6.4) and Proposition 6.1, we have

$$(6.16) \quad \mathcal{E}_f(U) - \mathcal{E}_f(U_0) \geq \frac{1}{2} (\mathcal{I}_1 + \mathcal{I}_2),$$

where we introduced the notation

$$\mathcal{I}_1 = \int_{\Omega} U u^2 dx - \left(\int_{\Omega} |u|^{2p/(p+1)} dx \right)^{(p+1)/p} \quad \text{and} \quad \mathcal{I}_2 = \|u - u_0\|_{W_0^{1,2}(\Omega)}^2.$$

We proceed to estimate \mathcal{I}_1 and \mathcal{I}_2 separately.

Estimate on \mathcal{I}_1 . For this we use the quantitative Hölder inequality (4.2) with

$$q = p + 1, \quad q' = \frac{p+1}{p}, \quad f = U^{-p/(p+1)}, \quad g = \frac{|u|^{2p/(p+1)} U^{p/(p+1)}}{\left(\int_{\Omega} U u^2 dx \right)^{p/(p+1)}}.$$

Thus we get

$$\begin{aligned} \left(\int_{\Omega} U u^2 dx \right)^{p/(p+1)} - \int_{\Omega} |u|^{2p/(p+1)} dx \\ \geq \frac{1}{2^p(p+1)} \left(\int_{\Omega} U u^2 dx \right)^{p/(p+1)} \left\| \frac{|u|^{2/(p+1)} U^{1/(p+1)}}{\left(\int_{\Omega} U u^2 dx \right)^{1/(p+1)}} - U^{-p/(p+1)} \right\|_{L^{p+1}(\Omega)}^{p+1}. \end{aligned}$$

By using (5.2) and (6.15) we have

$$\left(\int_{\Omega} U u^2 dx \right)^{\frac{p}{p+1}} \geq \left(\frac{c_2^2}{2} \right)^{\frac{p}{p+1}},$$

and by convexity of the function $t \mapsto t^{(p+1)/p}$, we have

$$\int_{\Omega} U u^2 dx - \left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx \right)^{\frac{p+1}{p}} \geq \frac{p+1}{p} \left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx \right)^{\frac{1}{p}} \left[\left(\int_{\Omega} U u^2 dx \right)^{\frac{p}{p+1}} - \int_{\Omega} |u|^{\frac{2p}{p+1}} dx \right].$$

Thus, for the moment we obtained

$$(6.17) \quad \mathcal{I}_1 \geq \frac{c_2^2}{p 2^{p+1}} \left\| \frac{|u|^{\frac{2}{p+1}} U^{\frac{1}{p+1}}}{\left(\int_{\Omega} u^2 U dx \right)^{\frac{1}{p+1}}} - \frac{1}{U^{\frac{p}{p+1}}} \right\|_{L^{p+1}(\Omega)}^{p+1},$$

where we used again (6.15) to estimate the norm of u from below. Observe that we have

$$\begin{aligned} \left\| \frac{|u|^{\frac{2}{p+1}} U^{\frac{1}{p+1}}}{\left(\int_{\Omega} u^2 U dx \right)^{\frac{1}{p+1}}} - \frac{1}{U^{\frac{p}{p+1}}} \right\|_{L^{p+1}(\Omega)}^{p+1} &= \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} u^2 U dx \right)^{\frac{1}{p+1}}} - \frac{1}{U} \right\|_{L^{p+1}(\Omega; U)}^{p+1} \\ &= \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} u^2 U dx \right)^{\frac{1}{p+1}}} - \frac{1}{U} \right\|_{L^2(\Omega; U)}^{\frac{p+1}{2} \cdot 2} \end{aligned}$$

then by applying (5.2) we get

$$\left\| \frac{|u|^{\frac{2}{p+1}} U^{\frac{1}{p+1}}}{\left(\int_{\Omega} u^2 U dx \right)^{\frac{1}{p+1}}} - \frac{1}{U^{\frac{p}{p+1}}} \right\|_{L^{p+1}(\Omega)}^{p+1} \geq \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} u^2 U dx \right)^{\frac{1}{p+1}}} - \frac{1}{U} \right\|_{L^p(\Omega)}^{p+1}.$$

We now use the triangle inequality and the convexity of $t \mapsto t^{p+1}$, so that

$$\begin{aligned} \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} u^2 U dx\right)^{\frac{1}{p+1}}} - \frac{1}{U} \right\|_{L^p(\Omega)}^{p+1} &\geq \frac{1}{2^p} \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}}} - \frac{1}{U} \right\|_{L^p(\Omega)}^{p+1} \\ &\quad - \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} u^2 U dx\right)^{\frac{1}{p+1}}} - \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}}} \right\|_{L^p(\Omega)}^{p+1}. \end{aligned}$$

The last term simply gives

$$\left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} u^2 U dx\right)^{\frac{1}{p+1}}} - \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}}} \right\|_{L^p(\Omega)}^{p+1} = \frac{\left| \left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}} - \left(\int_{\Omega} u^2 U dx\right)^{\frac{1}{p+1}} \right|^{p+1}}{\int_{\Omega} u^2 U dx}.$$

By keeping everything together, we have obtained

$$\begin{aligned} (6.18) \quad \mathcal{I}_1 &\geq \frac{c_2^2}{p 2^{2p+1}} \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}}} - \frac{1}{U} \right\|_{L^p(\Omega)}^{p+1} \\ &\quad - \frac{1}{p 2^p} \left| \left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}} - \left(\int_{\Omega} u^2 U dx\right)^{\frac{1}{p+1}} \right|^{p+1}, \end{aligned}$$

where as always we used (5.2) and (6.15) to estimate the $L^2(\Omega; U)$ norm of u . By using (6.16) and the convexity of $t \mapsto t^{p+1}$, we get

$$\begin{aligned} 2(\mathcal{E}_f(U) - \mathcal{E}_f(U_0)) &\geq \mathcal{I}_1 \geq (p+1) \int_{\Omega} |u|^{\frac{2p}{p+1}} dx \left[\left(\int_{\Omega} u^2 U dx\right)^{\frac{1}{p+1}} - \left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}} \right] \\ &\geq (p+1) \left(\frac{c_2^2}{2}\right)^{\frac{p}{p+1}} \left[\left(\int_{\Omega} u^2 U dx\right)^{\frac{1}{p+1}} - \left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}} \right] \end{aligned}$$

By using the latter, from (6.18) we can infer

$$\begin{aligned} \mathcal{I}_1 &\geq \frac{c_2^2}{p 2^{2p+1}} \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}}} - \frac{1}{U} \right\|_{L^p(\Omega)}^{p+1} \\ &\quad - \frac{2}{p(p+1)} \left(\frac{2}{(p+1)c_2^2}\right)^p (\mathcal{E}_f(U) - \mathcal{E}_f(U_0))^{p+1}. \end{aligned}$$

We now insert the previous estimate in (6.16), use that $\mathcal{E}_f(U) - \mathcal{E}_f(U_0) \leq 1$ and take the power $1/(p+1)$ on both sides. The resulting estimate is

$$(6.19) \quad \left(\mathcal{E}_f(U) - \mathcal{E}_f(U_0)\right)^{\frac{1}{p+1}} \geq c_6 \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx\right)^{\frac{1}{p}}} - \frac{1}{U} \right\|_{L^p(\Omega)},$$

where $c_6 > 0$ is the following constant depending only on p and c_2

$$(6.20) \quad c_6 := \left(1 + \frac{2}{p(p+1)} \left(\frac{2}{(p+1)c_2^2}\right)^p\right)^{-\frac{1}{p+1}} \left(\frac{c_2^2}{p4^{p+1}}\right)^{\frac{1}{p+1}}.$$

Estimate on \mathcal{I}_2 . Again by combining the triangle inequality and the convexity of $t \mapsto t^2$, we have

$$(6.21) \quad \begin{aligned} \mathcal{I}_2 = \|u - u_0\|_{W_0^{1,2}(\Omega)}^2 &= c_2^2 \left\| \frac{u}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} - \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right\|_{W_0^{1,2}(\Omega)}^2 \\ &\geq \frac{1}{2} c_2^2 \left\| \frac{u}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} - \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right\|_{W_0^{1,2}(\Omega)}^2 \\ &\quad - \left(\frac{\|u\|_{W_0^{1,2}(\Omega)}}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right)^2 \left(\|u\|_{L^{\frac{2p}{p+1}}(\Omega)} - \|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)} \right)^2. \end{aligned}$$

We also observe that, by recalling the energy estimate (2.3) and (6.15), we get

$$(6.22) \quad \left(\frac{\|u\|_{W_0^{1,2}(\Omega)}}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right)^2 \leq \frac{2}{c_2^2} \|f\|_{W^{-1,2}(\Omega)}^2.$$

In order to estimate the negative term on the right-hand side in (6.21), we can simply use Lemma 6.2. Indeed, we have

$$(6.23) \quad \begin{aligned} \left| \|u\|_{L^{\frac{2p}{p+1}}(\Omega)} - \|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)} \right| &\leq \left| \|u\|_{L^{\frac{2p}{p+1}}(\Omega)} - \|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)} \right| \frac{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)} + \|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \\ &= \frac{1}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \left| \|u\|_{L^{\frac{2p}{p+1}}(\Omega)}^2 - \|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}^2 \right| \\ &\leq \frac{1}{c_3 c_2} \sqrt{\mathcal{E}_f(U) - \mathcal{E}_f(U_0)} \end{aligned}$$

where we used (6.3) (recall that we are assuming (6.14)). By using (6.22) and (6.23) in (6.21), a further application of Sobolev inequality leads us to

$$(6.24) \quad \sqrt{\mathcal{E}_f(U) - \mathcal{E}_f(U_0)} \geq c_7 \left\| \frac{u}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} - \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right\|_{L^{2^*}(\Omega)},$$

where the constant $c_7 > 0$ depends only on N , f and c_2

$$(6.25) \quad c_7 := \sqrt{\frac{T_N}{c_3^2 c_2^4 + \|f\|_{W^{-1,2}(\Omega)}^2}} \frac{c_3 c_2^3}{2}.$$

Stability estimate for U . We now use Lemma 6.5 with the choices

$$r = \frac{2p}{p+1}, \quad s = 2^*,$$

and

$$(6.26) \quad g_0 = \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}}, \quad g = \frac{1}{2} \left[\frac{u}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} + \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right].$$

Thus we get

$$\|g\|_{L^{\frac{2p}{p+1}}(\Omega)} \geq 1 - c_4 c_2^{\frac{1-p}{p+1}} \|g - g_0\|_{L^{2^*}(\Omega)},$$

since

$$\frac{\| |g_0|^{r-1} \|_{L^{s'}(\Omega)}}{\|g_0\|_{L^r(\Omega)}^{r-1}} = \left\| |g_0|^{\frac{p-1}{p+1}} \right\|_{L^{(2^*)'(\Omega)}} = c_4 c_2^{\frac{1-p}{p+1}},$$

which is finite by hypothesis. By combining the previous with (6.24), we obtain¹

$$(6.27) \quad \|g\|_{L^{\frac{2p}{p+1}}(\Omega)} \geq 1 - \frac{c_4 c_2^{\frac{1-p}{p+1}}}{2 c_7} \sqrt{\mathcal{E}_f(U) - \mathcal{E}_f(U_0)}.$$

The previous estimate is crucial in order to estimate $g - g_0$ in $L^{2p/(p+1)}$. Indeed, by *Clarkson inequality*² and (6.27) we can infer

$$\|g - g_0\|_{L^{\frac{2p}{p+1}}(\Omega)}^{\frac{2p}{p-1}} \leq 1 - \|g\|_{L^{\frac{2p}{p+1}}(\Omega)}^{\frac{2p}{p-1}} \leq 1 - \left(1 - \frac{c_4 c_2^{\frac{1-p}{p+1}}}{2 c_7} \sqrt{\mathcal{E}_f(U) - \mathcal{E}_f(U_0)} \right)^{\frac{2p}{p-1}},$$

and thus

$$\|g - g_0\|_{L^{\frac{2p}{p+1}}(\Omega)}^{\frac{2p}{p-1}} \leq \frac{p}{p-1} \frac{c_4 c_2^{\frac{1-p}{p+1}}}{c_7} \sqrt{\mathcal{E}_f(U) - \mathcal{E}_f(U_0)},$$

thanks to the convexity of $t \mapsto t^{2p/(p-1)}$. We now go back to the definition (6.26) of g and g_0 , so that the previous finally gives

$$(6.28) \quad \left(\frac{1}{2} \right)^{\frac{2p}{p-1}} \left\| \frac{u}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} - \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right\|_{L^{\frac{2p}{p+1}}(\Omega)}^{\frac{2p}{p-1}} \leq \frac{p}{p-1} \frac{c_4 c_2^{\frac{1-p}{p+1}}}{c_7} \sqrt{\mathcal{E}_f(U) - \mathcal{E}_f(U_0)},$$

Recall that $2/(p+1) < 1$, thus the function $t \mapsto |t|^{2/(p+1)}$ is $2/(p+1)$ -Hölder continuous and we have

$$\left| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx \right)^{\frac{1}{p}}} - \frac{|u_0|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u_0|^{\frac{2p}{p+1}} dx \right)^{\frac{1}{p}}} \right|^p \leq \left| \frac{u}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} - \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right|^{\frac{2p}{p+1}}.$$

Thus from (6.28) we obtain

$$(6.29) \quad [\mathcal{E}_f(U) - \mathcal{E}_f(U_0)]^{\frac{p-1}{2p(p+1)}} \geq c_8 \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx \right)^{\frac{1}{p}}} - \frac{|u_0|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u_0|^{\frac{2p}{p+1}} dx \right)^{\frac{1}{p}}} \right\|_{L^p(\Omega)},$$

¹Up to further suppose that

$$\mathcal{E}_f(U) - \mathcal{E}_f(U_0) \leq \min \left\{ c_5, \left(\frac{2 c_7 c_2^{\frac{p-1}{p+1}}}{c_4} \right)^2 \right\},$$

we can assume that the right-hand side of (6.27) is positive.

²Let $1 < q \leq 2$ and $h_1, h_2 \in L^q(\Omega)$ be two functions with unit norm. Then we have

$$\left\| \frac{h_1 + h_2}{2} \right\|_{L^q(\Omega)}^{q'} + \left\| \frac{h_1 - h_2}{2} \right\|_{L^q(\Omega)}^{q'} \leq 1,$$

see [16, Theorem 2].

where the constant c_8 is given by

$$(6.30) \quad c_8 := \left(\frac{1}{4} \left(\frac{p-1}{p} \frac{c_7 c_2^{\frac{p-1}{p+1}}}{c_4} \right)^{\frac{p-1}{p}} \right)^{\frac{1}{p+1}}.$$

If we now use the relation (5.4) between U_0 and u_0 , the triangle inequality, (6.19) and (6.29) we get

$$\begin{aligned} \left\| \frac{1}{U} - \frac{1}{U_0} \right\|_{L^p(\Omega)} &\leq \left\| \frac{1}{U} - \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx \right)^{\frac{1}{p}}} \right\|_{L^p(\Omega)} + \left\| \frac{|u|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u|^{\frac{2p}{p+1}} dx \right)^{\frac{1}{p}}} - \frac{|u_0|^{\frac{2}{p+1}}}{\left(\int_{\Omega} |u_0|^{\frac{2p}{p+1}} dx \right)^{\frac{1}{p}}} \right\|_{L^p(\Omega)} \\ &\leq \left[\frac{\left(\mathcal{E}_f(U) - \mathcal{E}_f(U_0) \right)^{\frac{1}{p+1}}}{c_6} + \frac{\left(\mathcal{E}_f(U) - \mathcal{E}_f(U_0) \right)^{\frac{p-1}{2p(p+1)}}}{c_8} \right]. \end{aligned}$$

Observe that since $p > 1$, we have

$$\frac{p-1}{2p(p+1)} < \frac{1}{p+1},$$

thus it is now sufficient to use hypothesis (6.14) to get (6.12) for U .

Conclusion: stability estimate for V . We now go back to our potential V . By using (6.9), the previous step and (6.8), we get

$$\begin{aligned} \left\| \frac{1}{V} - \frac{1}{U_0} \right\|_{L^p(\Omega)} &\leq \left\| \frac{1}{U} - \frac{1}{U_0} \right\|_{L^p(\Omega)} + \frac{4}{c_2^2} \left(\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \right) \\ &\leq \frac{c_6 + c_8}{c_6 c_8} \left(\mathcal{E}_f(U) - \mathcal{E}_f(U_0) \right)^{\frac{p-1}{2p(p+1)}} + \frac{4}{c_2^2} \left(\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \right) \\ &\leq \left[\frac{c_6 + c_8}{c_6 c_8} + \frac{4}{c_2^2} \right] \left(\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \right)^{\frac{p-1}{2p(p+1)}}, \end{aligned}$$

where we also used that $\mathcal{E}_f(V) - \mathcal{E}_f(U_0) \leq 1$. This concludes the proof. \square

Some comments on the previous result are in order.

Remark 6.7 (Integrability assumption on u_0). We point out that

$$0 < (2^*)' \frac{p-1}{p+1} < 2, \quad \text{for every } p > 1,$$

thus the condition (6.11) of Theorem 6.6 is always satisfied if $|\Omega| < +\infty$. When $|\Omega| = +\infty$, this is still verified if f decreases sufficiently fast at infinity. For example, by appealing to Theorem A.1 in the Appendix this holds true for $f \in L^r(\Omega)$ with $r > N/2$ and

$$|f(x)| = O\left(\frac{1}{|x|^\alpha}\right), \quad \text{for } |x| \rightarrow \infty, \quad \alpha > \frac{N+2}{2}.$$

Observe that the condition on α is the minimal assumption for $|x|^{-\alpha}$ to be $(2^*)'$ -integrable at infinity.

Remark 6.8. A closer inspection of the previous proof informs us that a possible value for the constant σ_m in (6.12) is

$$\sigma_m = \min \left\{ \left(\frac{c_6 + c_8}{c_6 c_8} + \frac{4}{c_2^2} \right)^{-\frac{2p(p+1)}{p-1}}, c_5 4^{-\frac{p(p+1)}{p-1}}, \left(\frac{2c_7 c_2^{\frac{p-1}{p+1}}}{c_4} \right)^2 4^{-\frac{p(p+1)}{p-1}} \right\}.$$

where the constants c_5 , c_6 , c_7 and c_8 are defined in (6.13), (6.20), (6.25) and (6.30).

6.3. Stability of the state functions. By suitably combining some of the estimates we used so far, we also get a stability result for the energy functions in the natural space $W_0^{1,2}(\Omega) \cap L^{2p/(p+1)}(\Omega)$.

Proposition 6.9. *Under the hypotheses of Theorem 6.6, we have*

$$(6.31) \quad \left(\mathcal{E}_f(V) - \mathcal{E}_f(V_0) \right)^{\frac{p-1}{2p}} \geq c \left[\|u - u_0\|_{W_0^{1,2}(\Omega)}^2 + \|u - u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}^2 \right],$$

for some constant $c > 0$ depending on N, p, c_2, c_4 and $\|f\|_{W^{-1,2}(\Omega)}$.

Proof. We first observe that by (2.3) we have

$$\|u - u_0\|_{W_0^{1,2}(\Omega)} \leq \|u\|_{W_0^{1,2}(\Omega)} + \|u_0\|_{W_0^{1,2}(\Omega)} \leq 2\|f\|_{W^{-1,2}(\Omega)},$$

and by (5.3)

$$\|u - u_0\|_{L^{\frac{2p}{p+1}}(\Omega)} \leq \|u\|_{L^{\frac{2p}{p+1}}(\Omega)} + \|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)} \leq \sqrt{2}\|f\|_{W^{-1,2}(\Omega)}.$$

Thus we can assume without loss of generality that

$$\mathcal{E}_f(V) - \mathcal{E}_f(V_0) \leq 1,$$

otherwise the result is trivially true. From (6.4) and (6.2) we already know

$$\mathcal{E}_f(V) - \mathcal{E}_f(V_0) \geq \frac{1}{2} \|u - u_0\|_{W_0^{1,2}(\Omega)}^2$$

and by (6.28)

$$c_9 \left\| \frac{u}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} - \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right\|_{L^{\frac{4p}{p-1}}(\Omega)} \leq \mathcal{E}_f(V) - \mathcal{E}_f(V_0),$$

with

$$c_9 = \left(\frac{p-1}{p} \frac{c_7 c_2^{\frac{p-1}{p+1}}}{c_4} \right)^2 \left(\frac{1}{2} \right)^{\frac{4p}{p-1}}.$$

Moreover, by the triangle inequality it is not difficult to see that

$$\begin{aligned} \|u - u_0\|_{L^{\frac{2p}{p+1}}(\Omega)} &\leq \left| \|u\|_{L^{\frac{2p}{p+1}}(\Omega)} - \|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)} \right| \\ &\quad + c_2 \left\| \frac{u}{\|u\|_{L^{\frac{2p}{p+1}}(\Omega)}} - \frac{u_0}{\|u_0\|_{L^{\frac{2p}{p+1}}(\Omega)}} \right\|_{L^{\frac{2p}{p+1}}(\Omega)}. \end{aligned}$$

By combining these estimates and using (6.23), we get the desired conclusion. \square

Remark 6.10. By interpolation, it is easy to obtain a stability estimate like (6.31) in $L^r(\Omega)$ for every $2p/(p+1) < r < 2^*$ and in $W_0^{s,2}(\Omega)$ for every $0 < s < 1$.

APPENDIX A. SHARP DECAY ESTIMATES FOR NON AUTONOMOUS SCHRÖDINGER EQUATIONS

Given $1 < p < \infty$, we set for simplicity $q = 2p/(p+1)$ which is always between 1 and 2. In what follows we still denote by u_0 the unique minimizer of

$$J_{p,f}(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2} \left(\int_{\Omega} |u|^q dx \right)^{\frac{2}{q}} - \langle f, u \rangle.$$

The aim of this Appendix is to prove some decay properties for the optimal function u_0 , in the case $|\Omega| = +\infty$. We can confine ourselves to consider the case $\Omega = \mathbb{R}^N$.

Theorem A.1 (Properties of u_0). *Let $r > N/2$ and $f \in L^r(\mathbb{R}^N)$ be such that there exist $C, R > 0$ and $\alpha > (N+2)/2$ with*

$$|f(x)| \leq C |x|^{-\alpha}, \quad \text{for } |x| \geq R.$$

Then there exists $M = M(\|f\|_{L^r(\mathbb{R}^N)}, c_2, C, R, \alpha) > 0$ such that

$$(A.1) \quad |u_0(x)| \leq M, \quad x \in \mathbb{R}^N.$$

Moreover, if we denote by $w \in W_0^{1,2}(\mathbb{R}^N)$ the unique minimizer of

$$\mathcal{J}(u) = \frac{1}{2} \int_{\mathbb{R}^N} |\nabla u|^2 dx + \frac{c_2^{2-q}}{q} \int_{\mathbb{R}^N} |u|^q dx - \int_{\mathbb{R}^N} \frac{u}{(1+|x|^2)^{\alpha/2}} dx,$$

then there exists $T = T(M, C, R, \alpha) > 0$ such that

$$(A.2) \quad |u_0(x)| \leq T^{2/(2-q)} w\left(\frac{x}{T}\right), \quad x \in \mathbb{R}^N.$$

In particular, we get

$$(A.3) \quad |u_0(x)| \leq C' |x|^{-\alpha/(q-1)}, \quad \text{for } |x| \gg 1,$$

for some constant $C' > 0$.

Proof. We first observe that if $f \geq 0$, then the unique minimizer u_0 is positive since $J_{p,f}(|u|) \leq J_{p,f}(u)$. We also notice that it is not restrictive to prove the result for $f \geq 0$. Indeed, if f is not positive, by using the minimality of u_0 and the fact that $f \leq |f|$, it is not difficult to see that

$$|u_0| \leq \tilde{u}_0,$$

where \tilde{u}_0 is the unique minimizer of $J_{p,|f|}$. We thus assume $f \geq 0$ in what follows and divide the proof in three parts.

Boundedness of u_0 . The integrability of f already implies that $u_0 \in L_{loc}^\infty(\mathbb{R}^N)$, see [18, Chapter 7]. Also, since u_0 solves

$$-\Delta u_0 + c_2^{2-q} u_0^{q-1} = f,$$

it is the unique minimizer of the functional

$$\tilde{J}_{p,f}(u) = \frac{1}{2} \int_{\mathbb{R}^N} |\nabla u|^2 dx + \frac{c_2^{2-q}}{q} \int_{\mathbb{R}^N} |u|^q dx - \int_{\mathbb{R}^N} f u dx,$$

as well. Let $M > 0$, by testing the minimality of u_0 against $\varphi_M = \min\{u_0, M\}$, we get

$$\begin{aligned} \tilde{J}_{p,f}(\varphi_M) - \tilde{J}_{p,f}(u_0) &\leq \frac{c_2^{2-q}}{q} \int_{\{u_0 > M\}} M^q dx - \frac{c_2^{2-q}}{q} \int_{\{u_0 > M\}} u_0^q dx + \int_{\{u_0 > M\}} f(u_0 - M) dx \\ &= -\frac{c_2^{2-q}}{q} \int_{\{u_0 > M\}} (u_0^q - M^q) dx + \int_{\mathbb{R}^N} f(u_0 - M)_+ dx. \end{aligned}$$

We then observe that

$$u_0^q - M^q \geq q M^{q-1} (u_0 - M).$$

If we take

$$M = \max \left\{ \|u_0\|_{L^\infty(B_R)}, \left(c_2^{q-2} C R^{-\alpha} \right)^{1/(q-1)} \right\},$$

then we have

$$\{u_0 > M\} \subset \{x : |x| \geq R\},$$

so that

$$f(x) \leq C |x|^{-\alpha} \leq C R^{-\alpha} \leq c_2^{2-q} M^{q-1}, \quad \text{on } \{u_0 > M\}.$$

In conclusion, by using the choice of M and the decay of f , we get

$$\tilde{J}_{p,f}(\varphi_M) - \tilde{J}_{p,f}(u_0) \leq \int_{\mathbb{R}^N} \left(f - c_2^{2-q} M^{q-1} \right) (u_0 - M)_+ dx \leq 0.$$

By uniqueness of the minimizer of u_0 we get that $\varphi_M = u_0$ and thus $u_0 \leq M$ in \mathbb{R}^N .

Comparison. In order to prove the second assertion, we start observing that $\alpha > (N + 2)/2$ guarantees

$$h(x) = (1 + |x|^2)^{-\alpha/2} \in L^{(2^*)'}(\mathbb{R}^N) \subset W^{-1,2}(\mathbb{R}^N),$$

thus a function w minimizing \mathcal{J} exists, is unique and radially decreasing. Moreover, it solves

$$-\Delta w + c_2^{2-q} w^{q-1} = h, \quad \text{in } \mathbb{R}^N.$$

The rescaled function

$$w_t(x) = t^{2/(2-q)} w\left(\frac{x}{t}\right), \quad t > 0,$$

then solves

$$-\Delta w_t + c_2^{2-q} w_t^{q-1} = h_t, \quad \text{where} \quad h_t(x) = t^{2(q-1)/(2-q)} h\left(\frac{x}{t}\right).$$

Since by the first part of the proof u_0 is bounded and $w_{t_1} \geq w_{t_0}$ for $t_1 \geq t_0$, we can find a T_0 sufficiently large such that

$$(A.4) \quad w_t(x) \geq u_0(x), \quad \text{for } |x| \leq R \quad \text{and} \quad t \geq T_0.$$

In addition, if we define

$$T_1 = \max \left\{ R, (2^{\alpha/2} C R^{-\alpha})^{(2-q)/(2q-2)} \right\},$$

by hypothesis on f we get³

$$(A.5) \quad h_t(x) = \frac{t^{\frac{2q-2}{2-q} + \alpha}}{(t^2 + |x|^2)^{\alpha/2}} \geq \frac{t^{\frac{2q-2}{2-q}} R^\alpha}{2^{\alpha/2} |x|^\alpha} \geq \frac{C}{|x|^\alpha} \geq f(x), \quad \text{for } |x| > R \quad \text{and} \quad t \geq T_1.$$

We now define $T = \max\{T_0, T_1\}$ and test the minimality of u_0 against the function $\psi_T = \min\{u_0, w_T\}$. Thus we get

$$\begin{aligned} 0 \geq \tilde{J}_{p,f}(u_0) - \tilde{J}_{p,f}(\psi_T) &\geq \frac{1}{2} \int_{\{w_T < u_0\}} (|\nabla u_0|^2 - |\nabla w_T|^2) dx \\ &\quad + \frac{c_2^{2-q}}{q} \int_{\{w_T < u_0\}} (u_0^q - w_T^q) dx \\ &\quad - \int_{\{w_T < u_0\}} f(u_0 - w_T) dx. \end{aligned}$$

By using the convexity of the functions involved and the equation solved by w_T , we thus get

$$\begin{aligned} 0 \geq \int_{\mathbb{R}^N} \langle \nabla w_T, \nabla(u_0 - w_T)_+ \rangle dx + c_2^{2-q} \int_{\mathbb{R}^N} w_T^{q-1} (u_0 - w_T)_+ dx \\ - \int_{\mathbb{R}^N} f(u_0 - w_T)_+ dx = \int_{\mathbb{R}^N} (h_T - f)(u_0 - w_T)_+ dx. \end{aligned}$$

By combining (A.4) and (A.5), we have

$$h_T(x) \geq f(x), \quad \text{on } \{u_0 > w_T\},$$

thus we would obtain that ψ_T is a minimizer of $\tilde{J}_{p,f}$. By uniqueness, $\psi_T = u_0$ and thus (A.2) holds true.

Decay estimate for u_0 . Finally, the estimate (A.3) simply follows from (A.2) and Lemma A.3 below, applied to the rescaled function w_T . \square

³For the first inequality, we just use that for $x, t \geq R$ we have

$$\left(1 + \frac{|x|^2}{t^2}\right)^{\frac{\alpha}{2}} \leq \left(1 + \frac{|x|^2}{R^2}\right)^{\frac{\alpha}{2}} \leq 2^{\alpha/2} \frac{|x|^\alpha}{R^\alpha}.$$

Remark A.2. A different way to compare u_0 with a radial function and obtain (A.3) could be that of using symmetrization techniques. More precisely, one could look at the radial solution of the *symmetrized problem*

$$-\Delta v + c_2^{2-q} v^{q-1} = f^*,$$

where f^* denotes the *Schwarz rearrangement* of f (see [21, Chapter 2] for the relevant definition). There is a huge literature on results which permit to compare u_0^* and v (see for example [1, 26]), but the presence of the nonlinear term $c_2^{2-q} v^{q-1}$ complicates the task. An interesting result covering this case is contained in the recent paper [19] by Hamel and Russ, which however deals with the case of a *bounded domain* Ω . Since it is not clear whether this strategy could work or not, we decided to take a different path, which just uses the minimality of u_0 . We also refer to the related discussion in [27, Sections 5.2 and 5.3].

Lemma A.3 (Sharp decay estimate). *Let $N \geq 3$, $1 < q \leq 2$ and $\alpha > (N+2)/2$. Let us suppose that $u \in W_0^{1,2}(\mathbb{R}^N) \cap L^q(\mathbb{R}^N)$ is a smooth positive radial function verifying*

$$-\Delta u + a u^{q-1} \leq b |x|^{-\alpha}, \quad \text{for } |x| \geq R,$$

for some $a, R > 0$ and $b \geq 0$. Then there exists a constant $C > 0$ depending on $\alpha, q, R, a, \|u\|_{W_0^{1,2}(\mathbb{R}^N)}$ and $\|u\|_{L^q(\mathbb{R}^N)}$ such that

$$(A.6) \quad 0 \leq u(x) \leq C |x|^{-\alpha/(q-1)}, \quad \text{for } |x| \gg 1.$$

Proof. We divide the proof in two parts: in the first we prove that if

$$-\Delta u + a u^{q-1} \leq b |x|^{-\gamma}, \quad |x| \geq R,$$

for $\gamma > (N+2)/2$, then the following weaker decay estimate holds

$$(A.7) \quad 0 \leq u(x) \leq C_\varepsilon |x|^{-\gamma+\varepsilon}, \quad \text{for } |x| \gg 1,$$

for every $\varepsilon > 0$. Then we will get (A.6) by using a contradiction argument and a suitable maximum principle.

Part 1: weak decay. We first observe that Lemma A.5 below already implies that

$$u(\varrho) \lesssim \varrho^{-\beta_0}, \quad \text{where } \beta_0 = (N-1) \frac{2}{2+q}.$$

If $\gamma \leq \beta_0$, then (A.7) holds and there is nothing to prove. We can thus assume that $\gamma > \beta_0$. We are going to prove (A.7) by using a recursive argument, namely we will prove the following implication:

$$(A.8) \quad |u(\varrho)| \leq c \varrho^{-\beta}, \quad \text{for } \varrho > r_0 \text{ and } \beta \geq \beta_0 \implies |u(\varrho)| \leq c' \varrho^{-\frac{\gamma+\beta}{2}}, \quad \text{for } \varrho > r'_0.$$

By starting from $\beta = \beta_0$ and iterating a finite number of times (A.8), we will get the desired result. Indeed, observe that since $\gamma > \beta_0$, the sequence

$$\beta_{i+1} = \frac{\gamma + \beta_i}{2}$$

is monotone increasing and converges to γ .

To prove (A.8) we adapt a classical argument that can be found for example in [3, Lemma 2], but some care is needed in order to deal with the non-autonomous term. Also, for notational simplicity we give the proof just for $a = b = 1$. Then by hypothesis we have that u verifies

$$u''(\varrho) + \frac{N-1}{\varrho} u'(\varrho) - u(\varrho)^{q-1} \geq -\varrho^{-\gamma}, \quad \varrho \geq R.$$

The function $v(\varrho) = \varrho^{(N-1)/2} u(\varrho)$ verifies

$$v'' \geq \left[\frac{C_N}{\varrho^2} + (u+1)^{q-2} \right] v - \varrho^{\frac{N-1}{2}-\gamma}, \quad \varrho \geq R.$$

We then make the further substitution $w = v^2$

$$\frac{w''}{2} \geq (v')^2 + w \left[\frac{C_N}{\varrho^2} + (1+u)^{q-2} \right] - \varrho^{\frac{N-1}{2}-\gamma} \sqrt{w}, \quad \varrho \geq R.$$

As we already know that $u \rightarrow 0$ as $\varrho \rightarrow 0$, we have

$$\frac{C_N}{\varrho^2} + (1+u)^{q-2} \geq \frac{m}{2}, \quad \varrho \geq R,$$

where $m > 0$ is a suitable constant. Thus we obtain that w verifies

$$(A.9) \quad w'' - mw \geq -2\varrho^{\frac{N-1}{2}-\gamma} \sqrt{w}, \quad \varrho \geq R.$$

We then set

$$z(\varrho) = e^{-\sqrt{m}\varrho} (w'(\varrho) + \sqrt{m}w(\varrho)), \quad \varrho > R,$$

thus we get

$$z'(\varrho) = e^{-\sqrt{m}\varrho} (w''(\varrho) - mw(\varrho)) \geq -2e^{-\sqrt{m}\varrho} \varrho^{\frac{N-1}{2}-\gamma} \sqrt{w}, \quad \varrho > R.$$

where we used (A.9). In order to prove (A.8), we assume that

$$|u(\varrho)| \leq c\varrho^{-\beta}, \quad \varrho > r_0 \geq R,$$

for $\beta \geq \beta_0$, then by recalling that $w = \varrho^{N-1}u^2$ we get

$$\sqrt{w(\varrho)} \leq c\varrho^{\frac{N-1}{2}-\beta}, \quad \varrho > r_0.$$

Thus for z' we can infer

$$(A.10) \quad z'(\varrho) \geq -ce^{-\sqrt{m}\varrho} \varrho^{N-1-\gamma-\beta}, \quad \varrho > r_0,$$

Let us now take the (negative) function η defined by

$$\eta(\varrho) = -c \int_{\varrho}^{+\infty} \left[e^{-\sqrt{m}s} s^{N-1-\gamma-\beta} \right] ds, \quad \varrho > r_0.$$

Observe that η is strictly increasing and η goes to 0 as ϱ goes to ∞ . Then from (A.10) we get

$$z'(\varrho) \geq -ce^{-\sqrt{m}\varrho} \varrho^{N-1-\gamma-\beta} = -\eta'(\varrho), \quad \varrho > r_0,$$

that is $z + \eta$ is non decreasing on $(r_0, +\infty)$. Let us suppose that there exists $r_1 > r_0$ such that

$$z(r_1) + \eta(r_1) > 0,$$

then by monotonicity of $z + \eta$ we obtain

$$e^{\sqrt{m}\varrho} (z(\varrho) + \eta(\varrho)) \geq e^{\sqrt{m}\varrho} (z(r_1) + \eta(r_1)) > 0, \quad \varrho > r_1.$$

The previous gives a contradiction, since the right-hand side is not integrable on $(r_1, +\infty)$, while the left-hand side is. Indeed, observe that⁴

$$e^{\sqrt{m}\varrho} z(\varrho) = w'(\varrho) + \sqrt{m}w(\varrho) \in L^1((r_1, +\infty)),$$

and $e^{\sqrt{m}\varrho} \eta$ is integrable at infinity by construction, since

$$(A.11) \quad e^{\sqrt{m}\varrho} \eta(\varrho) \simeq \frac{\int_{\varrho}^{+\infty} \left[e^{-\sqrt{m}s} s^{N-1-\gamma-\beta} \right] ds}{e^{-\sqrt{m}\varrho}} \simeq \varrho^{N-1-\gamma-\beta},$$

and the latter is integrable on $(r_1, +\infty)$ thanks to the fact that $\gamma > (N+2)/2$ and $\beta \geq \beta_0$. From the previous argument, we get

$$z(\varrho) \leq -\eta(\varrho), \quad \varrho > r_0,$$

⁴We have that $w = \varrho^{N-1}u^2$ which is integrable near ∞ since $u \in W_0^{1,2}(\mathbb{R}^N) \cap L^q(\mathbb{R}^N) \subset L^2(\mathbb{R}^N)$. On the other hand, $w'(\varrho) = 2v(\varrho)v'(\varrho)$, with

$$v(\varrho) = \sqrt{w(\varrho)} \in L^2((r_1, +\infty)),$$

and

$$|v'(\varrho)| \lesssim \varrho^{\frac{N-3}{2}} |u(\varrho)| + \varrho^{\frac{N-1}{2}} |u'(\varrho)| \in L^2((r_1, +\infty)).$$

that is

$$\left(e^{\sqrt{m}\varrho} w(\varrho)\right)' = e^{\sqrt{m}\varrho} (w'(\varrho) + \sqrt{m} w(\varrho)) = e^{2\sqrt{m}\varrho} z(\varrho) \leq -e^{2\sqrt{m}\varrho} \eta(\varrho), \quad \varrho > r_0.$$

This in turn implies

$$0 \leq w(\varrho) \leq e^{-\sqrt{m}\varrho} \left[C - \int_R^\varrho e^{2\sqrt{m}s} \eta(s) ds \right], \quad \varrho > r_0.$$

By recalling the definition of η , we get

$$e^{-\sqrt{m}\varrho} \int_R^\varrho e^{2\sqrt{m}s} \eta(s) ds \simeq \frac{e^{2\sqrt{m}\varrho} \eta(\varrho)}{e^{\sqrt{m}\varrho}} \simeq e^{\sqrt{m}\varrho} \eta(\varrho) \simeq \varrho^{N-1-\gamma-\beta}, \quad \varrho \gg 1,$$

thanks to (A.11). This finally implies the following decay of w at infinity

$$w(\varrho) \leq C \varrho^{N-1-\gamma-\beta}, \quad \varrho \gg 1,$$

and by going back to u , we can finally infer

$$u(\varrho) = \sqrt{w} \varrho^{\frac{1-N}{2}} \leq \sqrt{C} \varrho^{\frac{1-N}{2}} \varrho^{\frac{N-1}{2}-\frac{\gamma}{2}-\frac{\beta}{2}} = \sqrt{C} \varrho^{-\frac{\gamma+\beta}{2}}, \quad \varrho \gg 1.$$

This concludes the proof of (A.8) and thus of (A.7), as already explained.

Part 2: sharp decay. We now prove (A.6). Let us assume by contradiction that u verifies

$$(A.12) \quad \lim_{|x| \rightarrow \infty} u(x) |x|^{\alpha/(q-1)} = +\infty.$$

This implies that for every $\varepsilon > 0$ there exists a radius R_ε such that

$$|x|^{-\alpha} < \varepsilon u(x)^{q-1}, \quad |x| \geq R_\varepsilon.$$

By taking $\varepsilon = a/2$, we thus get that u verifies

$$-\Delta u + \frac{a}{2} u^{q-1} \leq 0, \quad |x| \geq R_\varepsilon.$$

Thus by (A.7) we get that $u = o(|x|^{-\gamma})$ for every $\gamma > 0$, as $|x|$ goes to ∞ . This clearly contradicts (A.12), thus

$$0 \leq \liminf_{|x| \rightarrow \infty} u(x) |x|^{\alpha/(q-1)} < +\infty.$$

This implies that there exists a sequence $\{r_k\}_{k \in \mathbb{N}}$ of radii converging to ∞ and a constant $A > 0$ such that

$$u(x) \leq A |x|^{-\alpha/(q-1)}, \quad \text{for } k \in \mathbb{N} \text{ and } |x| = r_k.$$

We now take $\tilde{u}(x) = \tilde{A} |x|^{-\alpha/(q-1)}$, where $\tilde{A} \geq A$ is a constant large enough such that there exists a radius $\tilde{R} \gg 1$ for which

$$-\Delta \tilde{u} + a \tilde{u}^{q-1} \geq b |x|^{-\alpha}, \quad \text{for } |x| \geq \tilde{R}.$$

We take $k_0 = \min\{k : r_k \geq \tilde{R}\}$, then we claim that

$$(A.13) \quad u(x) \leq \tilde{u}(x), \quad \text{for } r_k \leq |x| \leq r_{k+1} \text{ and } k \geq k_0.$$

If (A.13) were not true, there would exist a radius r_k such that

$$\tilde{u}(y) - u(y) := \min_{r_k \leq |x| \leq r_{k+1}} (\tilde{u}(x) - u(x)) < 0.$$

Since by construction we have

$$0 \leq \tilde{u}(x) - u(x) \text{ for } |x| = r_k \quad \text{and} \quad 0 \leq \tilde{u}(x) - u(x) \text{ for } |x| = r_{k+1},$$

then y would be an interior minimum point of $\tilde{u} - u$. By using this and the differential inequalities verified by u and \tilde{u} , we would get

$$0 \leq \Delta \tilde{u}(y) - \Delta u(y) \leq a (\tilde{u}(y)^{q-1} - u(y)^{q-1}) < 0,$$

thanks to the strict monotonicity of $t \mapsto t^{q-1}$. This gives the desired contradiction, thus (A.13) holds true and the decay estimate on u is proved. \square

Remark A.4. Observe that the last part of the previous proof also shows that estimate (A.6) is the best possible.

In the previous proof we used the following result, which is essentially due to Strauss, see [24, Radial Lemma 1]. The statement is slightly more general (the original case corresponds to $q = 2$), the proof just relies upon Hölder inequality.

Lemma A.5 (Strauss lemma). *Let $N \geq 2$ and $u \in W_0^{1,2}(\mathbb{R}^N) \cap L^q(\mathbb{R}^N)$ be a radial function, where $0 < q < \infty$. Then we have*

$$(A.14) \quad |u(x)|^{2+q} \leq S_{q,N} |x|^{-2(N-1)} \left(\int_{\mathbb{R}^N} |\nabla u|^2 dx \right) \left(\int_{\mathbb{R}^N} |u|^q dx \right), \quad x \in \mathbb{R}^N,$$

where

$$S_{q,N} = \left(\frac{1}{N \omega_N} \frac{2+q}{2} \right)^2$$

and ω_N is the measure of the N -dimensional ball of radius 1.

Proof. Let $u \in C_0^\infty(\mathbb{R}^N)$ be a radial function, for every $p > 1$ we have (with a small abuse of notation)

$$|u(\varrho)|^p = - \int_{\varrho}^{+\infty} \frac{d}{dt} |u(t)|^p dt = -p \int_{\varrho}^{+\infty} u'(t) |u(t)|^{p-2} u(t) dt, \quad \varrho > 0.$$

By taking $p = (2+q)/2$, we thus get

$$\begin{aligned} |u(\varrho)|^{(2+q)/2} &\leq \frac{2+q}{2} \int_{\varrho}^{+\infty} |u'(t)| |u(t)|^{q/2} dt \\ &\leq \frac{2+q}{2} \varrho^{1-N} \int_{\varrho}^{+\infty} \left(|u'(t)| t^{(N-1)/2} \right) \left(|u(t)|^{q/2} t^{(N-1)/2} \right) dt, \quad \varrho > 0. \end{aligned}$$

We now use Hölder inequality, then for every $\varrho > 0$

$$(A.15) \quad |u(\varrho)|^{(2+q)/2} \leq \frac{2+q}{2} \varrho^{1-N} \left(\int_{\varrho}^{+\infty} |u'(t)|^2 t^{N-1} dt \right)^{1/2} \left(\int_{\varrho}^{+\infty} |u(t)|^q t^{N-1} dt \right)^{1/2}.$$

By noticing that for a radial function

$$\left(\int_{\mathbb{R}^N} |\nabla u|^2 dx \right)^{1/2} = \sqrt{N \omega_N} \left(\int_0^{+\infty} |u'(t)|^2 t^{N-1} dt \right)^{1/2},$$

and

$$\left(\int_{\mathbb{R}^N} |u|^q dx \right)^{1/q} = (N \omega_N)^{1/q} \left(\int_0^{+\infty} |u(t)|^q t^{N-1} dt \right)^{1/q}$$

from (A.15) we finally get (A.14) for smooth functions. The inequality for $u \in W_0^{1,2}(\mathbb{R}^N) \cap L^q(\mathbb{R}^N)$ is obtained by a standard density argument. \square

Acknowledgements. We warmly thank François Hamel for some valuable suggestions on the proof of Lemma A.3. The work of the second author is part of the project 2010A2TFX2 “Calcolo delle Variazioni” funded by the Italian Ministry of Research and University. Part of this work has been done during some visits of the first author to Pisa. The Departement of Mathematics of the University of Pisa and its facilities are kindly acknowledged.

REFERENCES

- [1] A. Alvino, P.-L. Lions, G. Trombetti, Comparison results for elliptic and parabolic equations via Schwarz symmetrization, *Ann. Inst. Henri Poincaré*, **7** (1990), 37–65.
- [2] M. S. Ashbaugh, E. M. Harrell, Maximal and minimal eigenvalues and their associated nonlinear equations, *J. Math. Phys.*, **28** (1987), 1770–1786.
- [3] H. Berestycki, P.-L. Lions, Nonlinear scalar field equations. I. Existence of a ground state, *Arch. Rational Mech. Anal.*, **82** (1983), 313–345.
- [4] G. Bianchi, H. Egnell, A note on the Sobolev inequality, *J. Funct. Anal.*, **100** (1991), 18–24.

- [5] L. Brasco, G. De Philippis, B. Velichkov, Faber-Krahn inequalities in sharp quantitative form, preprint (2013), available at <http://arxiv.org/abs/1306.0392>
- [6] L. Brasco, G. De Philippis, B. Ruffini, Spectral optimization for the Stekloff-Laplacian: the stability issue, *J. Funct. Anal.*, **262** (2012), 4675–4710.
- [7] L. Brasco, A. Pratelli, Sharp stability of some spectral inequalities, *Geom. Funct. Anal.*, **22** (2012), 107–135.
- [8] D. Bucur, G. Buttazzo, B. Velichkov, Spectral optimization problems for potentials and measures, preprint (2013), available at <http://cvgmt.sns.it/paper/2245/>
- [9] G. Buttazzo, Spectral optimization problems, *Rev. Mat. Complut.*, **24** (2011), 277–322.
- [10] G. Buttazzo, A. Gerolin, B. Ruffini, B. Velichkov, Optimal Potentials for Schrödinger Operators, *J. Ecole Polytechnique*, **1** (2014), 71–100.
- [11] E. A. Carlen, A. Figalli, Stability for a GNS inequality and the Log-HLS inequality, with application to the critical mass Keller-Segel equation, *Duke Math. J.*, **162** (2013), 579–625.
- [12] E. A. Carlen, R. L. Frank, E. H. Lieb, Stability estimates for the lowest eigenvalue of a Schrödinger operator, *Geom. Funct. Anal.*, (to appear), available at <http://arxiv.org/abs/1301.5032>.
- [13] M. Cicalese, G. P. Leonardi, A Selection Principle for the Sharp Quantitative Isoperimetric Inequality, *Arch. Rational Mech. Anal.*, **206** (2012), 617–643.
- [14] A. Cianchi, L. Esposito, N. Fusco, C. Trombetti, A quantitative Pólya-Szegő principle, *J. Reine Angew. Math.*, **614** (2008), 153–189.
- [15] A. Cianchi, N. Fusco, F. Maggi, A. Pratelli, The sharp Sobolev inequality in quantitative form, *J. Eur. Math. Soc.*, **11** (2009), 1105–1139.
- [16] J. A. Clarkson, Uniformly convex spaces, *Trans. Am. Math. Soc.*, **40** (1936), 396–414.
- [17] N. Fusco, F. Maggi, A. Pratelli, The sharp quantitative isoperimetric inequality, *Ann. of Math.*, **168** (2008), 941–980.
- [18] E. Giusti, *Metodi diretti nel calcolo delle variazioni*. (Italian) [Direct methods in the calculus of variations], Unione Matematica Italiana, Bologna, 1994.
- [19] F. Hamel, E. Russ, Comparison results for semilinear elliptic equations using a new symmetrization method, preprint (2014), available at <http://arxiv.org/abs/1401.1726>
- [20] W. Hansen, N. Nadirashvili, Isoperimetric inequalities in potential theory, *Potential Anal.*, **3** (1994), 1–14.
- [21] A. Henrot, *Extremum Problems for Eigenvalues of Elliptic Operators*, *Frontiers in Mathematics*, Birkhäuser Verlag, Basel (2006).
- [22] J. B. Keller, Lower Bounds and Isoperimetric Inequalities for Eigenvalues of the Schrödinger Equation, *J. Math. Phys.*, **2** (1961), 262–266.
- [23] A. Melas, The stability of some eigenvalue estimates, *J. Differential Geom.*, **36** (1992), 19–33.
- [24] W. A. Strauss, Existence of Solitary Waves in Higher Dimensions, *Commun. Math. Phys.*, **55** (1977), 149–162.
- [25] G. Talenti, Elliptic equations and rearrangements, *Ann. Sc. Norm. Sup. Pisa*, **3** (1976), 697–718.
- [26] G. Talenti, Best constant in Sobolev inequality, *Ann. Mat. Pura Appl.*, **110** (1976), 353–372.
- [27] J. L. Vazquez, Symmetrization and mass comparison for degenerate nonlinear parabolic and related elliptic equations, *Adv. Nonlinear Stud.*, **5** (2005), 87–131.

AIX-MARSEILLE UNIVERSITÉ, CNRS, CENTRALE MARSEILLE, I2M, UMR 7373, 13453 MARSEILLE, FRANCE
E-mail address: lorenzo.brasco@univ-amu.fr

DIPARTIMENTO DI MATEMATICA, UNIVERSITÀ DI PISA, LARGO B. PONTECORVO 5, 56126 PISA, ITALY
E-mail address: buttazzo@dm.unipi.it